

Topic 7

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# Two- and Three-Dimensional Solid Elements; Plane Stress, Plane Strain, and Axisymmetric Conditions

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**Contents:**

- Isoparametric interpolations of coordinates and displacements
- Consistency between coordinate and displacement interpolations
- Meaning of these interpolations in large displacement analysis, motion of a material particle
- Evaluation of required derivatives
- The Jacobian transformations
- Details of strain-displacement matrices for total and updated Lagrangian formulations
- Example of 4-node two-dimensional element, details of matrices used

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**Textbook:**

Sections 6.3.2, 6.3.3

**Example:**

6.17

- FINITE ELEMENTS CAN IN GENERAL BE CATEGORIZED AS
  - CONTINUUM ELEMENTS (SOLID)
  - STRUCTURAL ELEMENTS

IN THIS LECTURE

- WE CONSIDER THE 2-D CONTINUUM ISOPARAMETRIC ELEMENTS
- THESE ELEMENTS ARE USED VERY WIDELY

- THE ELEMENTS ARE VERY GENERAL ELEMENTS FOR GEOMETRIC AND MATERIAL NONLINEAR CONDITIONS
- WE ALSO POINT OUT HOW GENERAL 3-D ELEMENTS ARE CALCULATED USING THE SAME PROCEDURES

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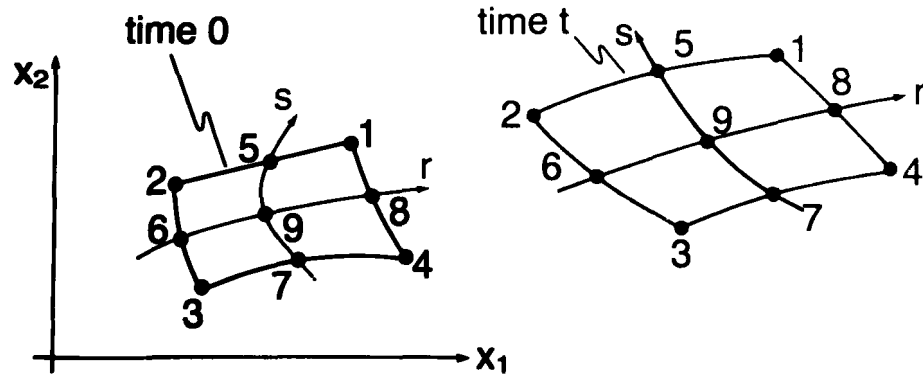
### TWO- AND THREE-DIMENSIONAL SOLID ELEMENTS

- Two-dimensional elements comprise
  - plane stress and plane strain elements
  - axisymmetric elements
- The derivations used for the two-dimensional elements can be easily extended to the derivation of three-dimensional elements.

Hence we concentrate our discussion now first on the two-dimensional elements.

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### TWO-DIMENSIONAL AXISYMMETRIC, PLANE STRAIN AND PLANE STRESS ELEMENTS



Because the elements are isoparametric,

$${}^0x_1 = \sum_{k=1}^N h_k {}^0x_1^k, \quad {}^0x_2 = \sum_{k=1}^N h_k {}^0x_2^k$$

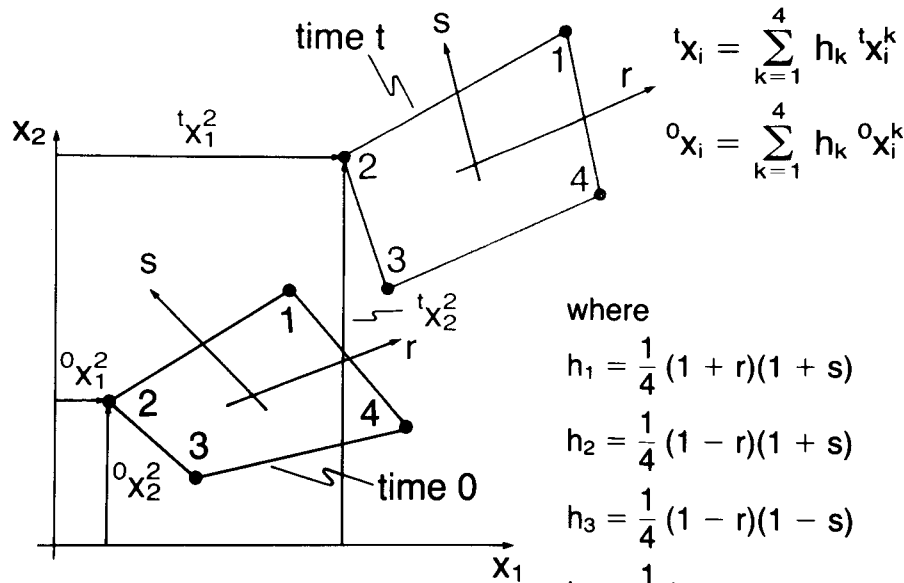
and

$${}^tx_1 = \sum_{k=1}^N h_k {}^tx_1^k, \quad {}^tx_2 = \sum_{k=1}^N h_k {}^tx_2^k$$

where the  $h_k$ 's are the isoparametric interpolation functions.

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Example: A four-node element



$${}^tx_i = \sum_{k=1}^4 h_k {}^tx_i^k$$

$${}^0x_i = \sum_{k=1}^4 h_k {}^0x_i^k$$

where

$$h_1 = \frac{1}{4} (1 + r)(1 + s)$$

$$h_2 = \frac{1}{4} (1 - r)(1 + s)$$

$$h_3 = \frac{1}{4} (1 - r)(1 - s)$$

$$h_4 = \frac{1}{4} (1 + r)(1 - s)$$

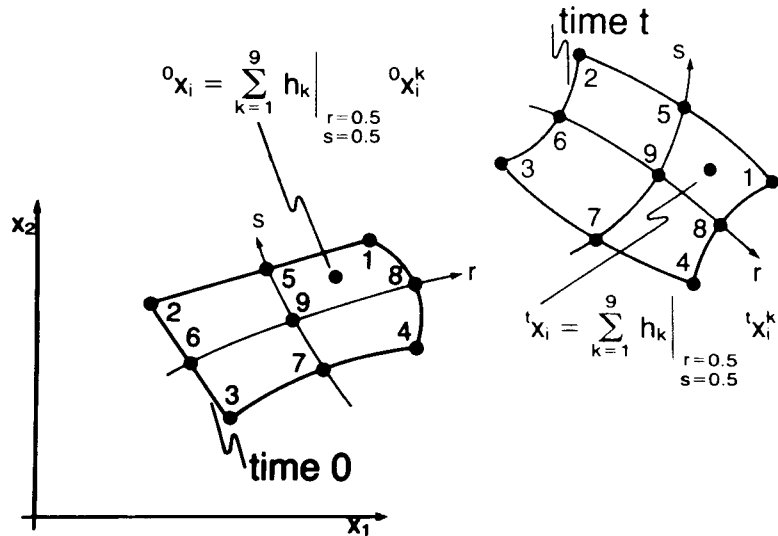
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### Example: Motion of a material particle

Consider the material particle at  $r = 0.5, s = 0.5$ :

Important: The isoparametric coordinates of a material particle never change



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A major advantage of the isoparametric finite element discretization is that we may directly write

$${}^t u_1 = \sum_{k=1}^N h_k {}^t u_1^k \quad , \quad {}^t u_2 = \sum_{k=1}^N h_k {}^t u_2^k$$

and

$$u_1 = \sum_{k=1}^N h_k u_1^k \quad , \quad u_2 = \sum_{k=1}^N h_k u_2^k$$

This is easily shown: for example,

$${}^t x_i = \sum_{k=1}^N h_k {}^t x_i^k$$

$${}^o x_i = \sum_{k=1}^N h_k {}^o x_i^k$$

Subtracting the second equation from the first equation gives

$$\underbrace{{}^t x_i - {}^o x_i}_{{}^t u_i} = \sum_{k=1}^N h_k \underbrace{({}^t x_i^k - {}^o x_i^k)}_{{}^t u_i^k}$$

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The element matrices require the following derivatives:

$${}^t u_{i,j} = \frac{\partial {}^t u_i}{\partial {}^o x_j} = \sum_{k=1}^N \left( \frac{\partial h_k}{\partial {}^o x_j} \right) {}^t u_i^k$$

$${}^o u_{i,j} = \frac{\partial u_i}{\partial {}^o x_j} = \sum_{k=1}^N \left( \frac{\partial h_k}{\partial {}^o x_j} \right) u_i^k$$

$${}^t u_{i,j} = \frac{\partial u_i}{\partial {}^t x_j} = \sum_{k=1}^N \left( \frac{\partial h_k}{\partial {}^t x_j} \right) u_i^k$$

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These derivatives are evaluated using a Jacobian transformation (the chain rule):

$$\frac{\partial h_k}{\partial r} = \frac{\partial h_k}{\partial^0 x_1} \frac{\partial^0 x_1}{\partial r} + \frac{\partial h_k}{\partial^0 x_2} \frac{\partial^0 x_2}{\partial r}$$

$$\frac{\partial h_k}{\partial s} = \frac{\partial h_k}{\partial^0 x_1} \frac{\partial^0 x_1}{\partial s} + \frac{\partial h_k}{\partial^0 x_2} \frac{\partial^0 x_2}{\partial s}$$

In matrix form,

$$\begin{bmatrix} \frac{\partial h_k}{\partial r} \\ \frac{\partial h_k}{\partial s} \end{bmatrix} = \underbrace{\begin{bmatrix} \frac{\partial^0 x_1}{\partial r} & \frac{\partial^0 x_2}{\partial r} \\ \frac{\partial^0 x_1}{\partial s} & \frac{\partial^0 x_2}{\partial s} \end{bmatrix}}_{{}^0 J} \begin{bmatrix} \frac{\partial h_k}{\partial^0 x_1} \\ \frac{\partial h_k}{\partial^0 x_2} \end{bmatrix}$$

REQUIRED DERIVATIVES

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The required derivatives are computed using a matrix inversion:

$$\begin{bmatrix} \frac{\partial h_k}{\partial^0 x_1} \\ \frac{\partial h_k}{\partial^0 x_2} \end{bmatrix} = {}^0 J^{-1} \begin{bmatrix} \frac{\partial h_k}{\partial r} \\ \frac{\partial h_k}{\partial s} \end{bmatrix}$$

The entries in  ${}^0 J$  are computed using the interpolation functions. For example,

$$\frac{\partial^0 x_1}{\partial r} = \sum_{k=1}^N \frac{\partial h_k}{\partial r} {}^0 x_1^k$$

The derivatives taken with respect to the configuration at time  $t$  can also be evaluated using a Jacobian transformation.

$$\begin{bmatrix} \frac{\partial h_k}{\partial r} \\ \frac{\partial h_k}{\partial s} \end{bmatrix} = \begin{bmatrix} \frac{\partial^t x_1}{\partial r} & \frac{\partial^t x_2}{\partial r} \\ \frac{\partial^t x_1}{\partial s} & \frac{\partial^t x_2}{\partial s} \end{bmatrix} \begin{bmatrix} \frac{\partial h_k}{\partial^t x_1} \\ \frac{\partial h_k}{\partial^t x_2} \end{bmatrix}$$

$\underline{J}^t$

$$\begin{bmatrix} \frac{\partial h_k}{\partial^t x_1} \\ \frac{\partial h_k}{\partial^t x_2} \end{bmatrix} = \underline{J}^{t-1} \begin{bmatrix} \frac{\partial h_k}{\partial r} \\ \frac{\partial h_k}{\partial s} \end{bmatrix}$$

$\sum_{k=1}^N \frac{\partial h_k}{\partial s} \frac{\partial^t x_2^k}{\partial s}$

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We can now compute the required element matrices for the total Lagrangian formulation:

Element Matrix	Matrices Required
$\underline{0}^t \underline{K}_L$	$\underline{0}^t \underline{C}$ , $\underline{0}^t \underline{B}_L$
$\underline{0}^t \underline{K}_{NL}$	$\underline{0}^t \underline{S}$ , $\underline{0}^t \underline{B}_{NL}$
$\underline{0}^t \underline{F}$	$\underline{0}^t \underline{\hat{S}}$ , $\underline{0}^t \underline{B}_L$

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We define  ${}^0\underline{C}$  so that

$$\begin{bmatrix} {}^0S_{11} \\ {}^0S_{22} \\ {}^0S_{12} \\ {}^0S_{33} \end{bmatrix} = {}^0\underline{C} \begin{bmatrix} {}^0e_{11} \\ {}^0e_{22} \\ 2 {}^0e_{12} \\ {}^0e_{33} \end{bmatrix}$$

analogous to  
 ${}^0S_{ij} = {}^0C_{ijrs} {}^0e_{rs}$

For example, we may choose  
(axisymmetric analysis),

$${}^0\underline{C} = \frac{E(1-\nu)}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1 & \frac{\nu}{1-\nu} & 0 & \frac{\nu}{1-\nu} \\ \frac{\nu}{1-\nu} & 1 & 0 & \frac{\nu}{1-\nu} \\ 0 & 0 & \frac{1-2\nu}{2(1-\nu)} & 0 \\ \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & 0 & 1 \end{bmatrix}$$

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We note that, in two-dimensional  
analysis,

$$\begin{aligned} {}^0e_{11} &= {}^0u_{1,1} + \underbrace{{}^t u_{1,1} \quad {}^0u_{1,1} + {}^t u_{2,1} \quad {}^0u_{2,1}} \\ {}^0e_{22} &= {}^0u_{2,2} + \underbrace{{}^t u_{1,2} \quad {}^0u_{1,2} + {}^t u_{2,2} \quad {}^0u_{2,2}} \\ 2 {}^0e_{12} &= ({}^0u_{1,2} + {}^0u_{2,1}) + \underbrace{({}^t u_{1,1} \quad {}^0u_{1,2} \\ &+ {}^t u_{2,1} \quad {}^0u_{2,2} + {}^t u_{1,2} \quad {}^0u_{1,1} + {}^t u_{2,2} \quad {}^0u_{2,1})} \\ {}^0e_{33} &= \frac{u_1}{{}^0x_1} + \underbrace{\begin{pmatrix} {}^t u_1 \\ {}^0x_1 \end{pmatrix} \frac{u_1}{{}^0x_1}} \end{aligned}$$

INITIAL DISPLACEMENT  
EFFECT

and

$${}^0\eta_{11} = \frac{1}{2} (({}^0u_{1,1})^2 + ({}^0u_{2,1})^2)$$

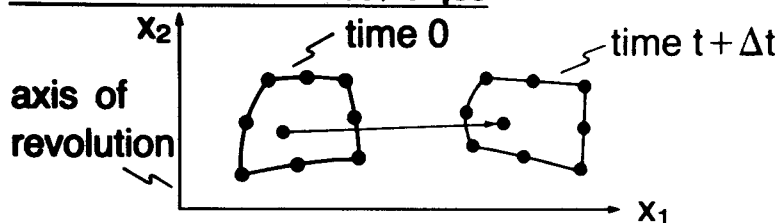
$${}^0\eta_{22} = \frac{1}{2} (({}^0u_{1,2})^2 + ({}^0u_{2,2})^2)$$

$${}^0\eta_{12} = {}^0\eta_{21} = \frac{1}{2} ({}^0u_{1,1} {}^0u_{1,2} + {}^0u_{2,1} {}^0u_{2,2})$$

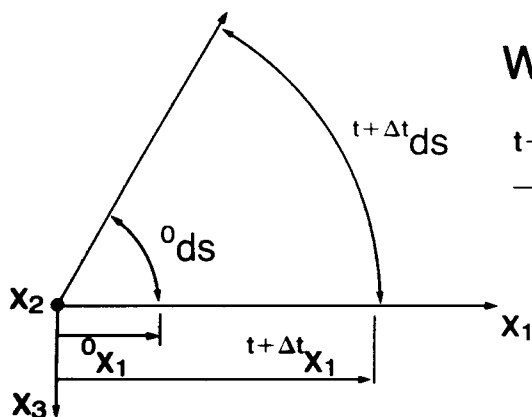
$${}^0\eta_{33} = \frac{1}{2} \left( \frac{u_1}{{}^0x_1} \right)^2$$

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**Derivation of  ${}^0e_{33}$ ,  ${}^0\eta_{33}$ :**



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We see that

$$\frac{{}^{t+\Delta t}ds}{{}^0ds} = \frac{{}^{t+\Delta t}x_1}{{}^0x_1}$$

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Hence

$$\begin{aligned}
 {}^{t+\Delta t}{}_0\epsilon_{33} &= \frac{1}{2} \left[ \left( \frac{{}^{t+\Delta t}ds}{{}_0ds} \right)^2 - 1 \right] \\
 &= \frac{1}{2} \left[ \left( \frac{{}^{t+\Delta t}x_1}{{}_0x_1} \right)^2 - 1 \right] \\
 &= \frac{1}{2} \left[ \left( \frac{{}_0x_1 + {}^t u_1 + u_1}{{}_0x_1} \right)^2 - 1 \right] \\
 &\vdots \\
 &= \underbrace{\left( \frac{{}^t u_1}{{}_0x_1} + \frac{1}{2} \left( \frac{{}^t u_1}{{}_0x_1} \right)^2 \right)}_{{}_0\epsilon_{33}} \\
 &\quad + \underbrace{\left( \frac{u_1}{{}_0x_1} + \left( \frac{{}^t u_1}{{}_0x_1} \right) \frac{u_1}{{}_0x_1} + \frac{1}{2} \left( \frac{u_1}{{}_0x_1} \right)^2 \right)}_{{}_0\epsilon_{33}}
 \end{aligned}$$

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We construct  ${}^t\underline{B}_L$  so that

$$\begin{bmatrix} {}_0e_{11} \\ {}_0e_{22} \\ 2{}_0e_{12} \\ {}_0e_{33} \end{bmatrix} = \underline{0e} = \underbrace{({}^t\underline{B}_{L0} + \underbrace{{}^t\underline{B}_{L1}}_{\text{contains initial displacement effect}})}_{{}^t\underline{B}_L} \hat{u}$$

${}_0e_{33}$  is only included for axisymmetric analysis

Entries in  ${}^t\mathbf{B}_{L0}$ :

		node k			
		$u_1^k$	$u_2^k$		
...	$\frac{\partial h_{k,1}}{\partial u_1^k}$	0	0	$\frac{\partial h_{k,2}}{\partial u_1^k}$	...
$\dots$	0	$\frac{\partial h_{k,2}}{\partial u_2^k}$	$\frac{\partial h_{k,1}}{\partial u_2^k}$	$\dots$	$\dots$
$\dots$	$\frac{\partial h_{k,2}}{\partial u_1^k}$	$\frac{\partial h_{k,1}}{\partial u_2^k}$	0	$\dots$	$\dots$
$\dots$	$\frac{h_k}{\sigma_{X_1}}$	0	0	$\dots$	$\dots$

included only for  
axisymmetric analysis

This is similar in form to the  $\mathbf{B}$  matrix used in linear analysis.

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Entries in  ${}^t\mathbf{B}_{L1}$ : node k

		node k			
		$u_1^k$	$u_2^k$		
$\dots$	$\frac{\partial u_{1,1}}{\partial u_1^k} \frac{\partial h_{k,1}}{\partial u_1^k}$	$\frac{\partial u_{1,2}}{\partial u_1^k} \frac{\partial h_{k,2}}{\partial u_1^k}$	$\frac{\partial u_{2,1}}{\partial u_2^k} \frac{\partial h_{k,1}}{\partial u_2^k}$	$\frac{\partial u_{2,2}}{\partial u_2^k} \frac{\partial h_{k,2}}{\partial u_2^k}$	$\dots$
$\dots$	$\frac{\partial u_{1,1}}{\partial u_1^k} \frac{\partial h_{k,2}}{\partial u_1^k}$	$+$ $\frac{\partial u_{1,2}}{\partial u_1^k} \frac{\partial h_{k,1}}{\partial u_1^k}$	$\frac{\partial u_{2,1}}{\partial u_2^k} \frac{\partial h_{k,2}}{\partial u_2^k}$	$+$ $\frac{\partial u_{2,2}}{\partial u_2^k} \frac{\partial h_{k,1}}{\partial u_2^k}$	$\dots$
$\dots$	$\frac{t_{u_1}}{\sigma_{X_1}} \frac{h_k}{\sigma_{X_1}}$	$0$	$0$	$0$	$\dots$

The initial displacement effect is contained in the terms  $\frac{t_{u_i}}{\sigma_{X_1}}$ ,  $\frac{t_{u_1}}{\sigma_{X_1}}$ .

included only  
for axisymmetric  
analysis

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Transparency  
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We construct  ${}^t\underline{B}_{NL}$  and  ${}^t\underline{S}$  so that

$$\delta \underline{\hat{u}}^T {}^t\underline{B}_{NL}^T {}^t\underline{S} {}^t\underline{B}_{NL} \underline{\hat{u}} = {}^t\underline{S}_{ij} \delta \eta_{ij}$$

Entries in  ${}^t\underline{S}$ :

${}^tS_{11}$	${}^tS_{12}$	0	0	0
${}^tS_{21}$	${}^tS_{22}$	0	0	0
0	0	${}^tS_{11}$	${}^tS_{12}$	0
0	0	${}^tS_{21}$	${}^tS_{22}$	0
0	0	0	0	${}^tS_{33}$

included only  
for axisymmetric  
analysis

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Entries in  ${}^t\underline{B}_{NL}$ :

		node k			
		$u_1^k$	$u_2^k$		
...	${}^0h_{k,1}$	0	0	...	$\vdots$
	${}^0h_{k,2}$	0	0		$u_1^k$
	0	0	${}^0h_{k,1}$		$u_2^k$
	0	0	${}^0h_{k,2}$		$\vdots$
$\swarrow$	$h_k/{}^0x_1$	0	0		

node k

included only for  
axisymmetric  
analysis

${}^t\hat{\underline{S}}$  is constructed so that

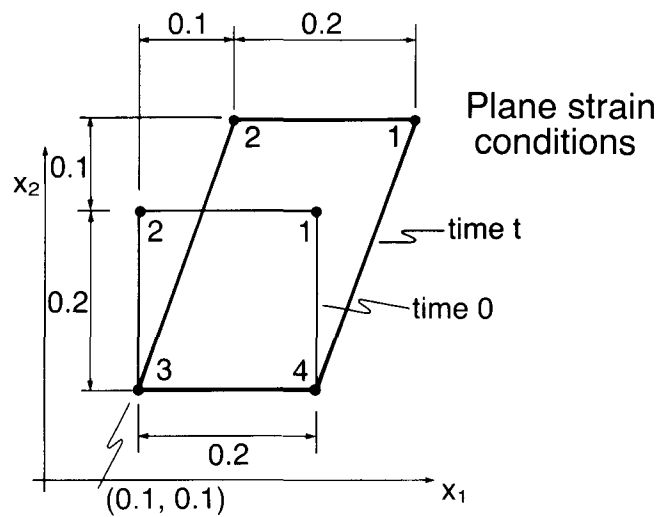
$$\delta \underline{\hat{u}}^T {}^t\underline{B}_L^T {}^t\hat{\underline{S}} = {}^tS_{ij} \delta e_{ij}$$

Entries in  ${}^t\hat{\underline{S}}$ :

$$\begin{bmatrix} {}^tS_{11} \\ {}^tS_{22} \\ \frac{{}^tS_{12}}{2} \\ {}^tS_{33} \end{bmatrix} \left\{ \begin{array}{l} \text{--- included only for} \\ \text{axisymmetric analysis} \end{array} \right.$$

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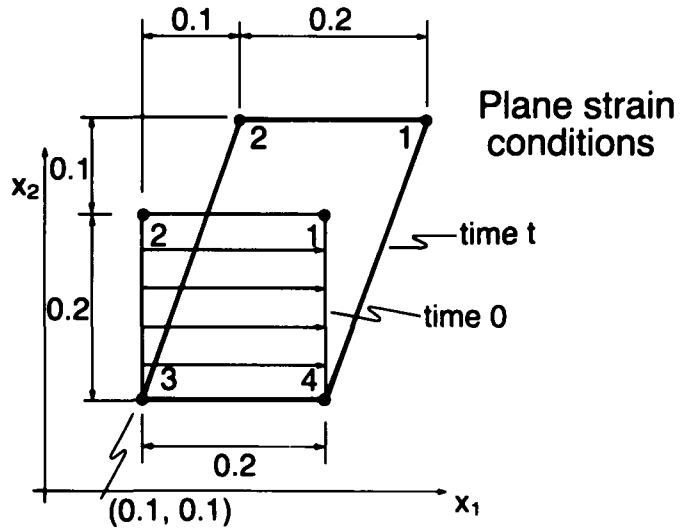
Example: Calculation of  ${}^t\underline{B}_L, {}^t\underline{B}_{NL}$



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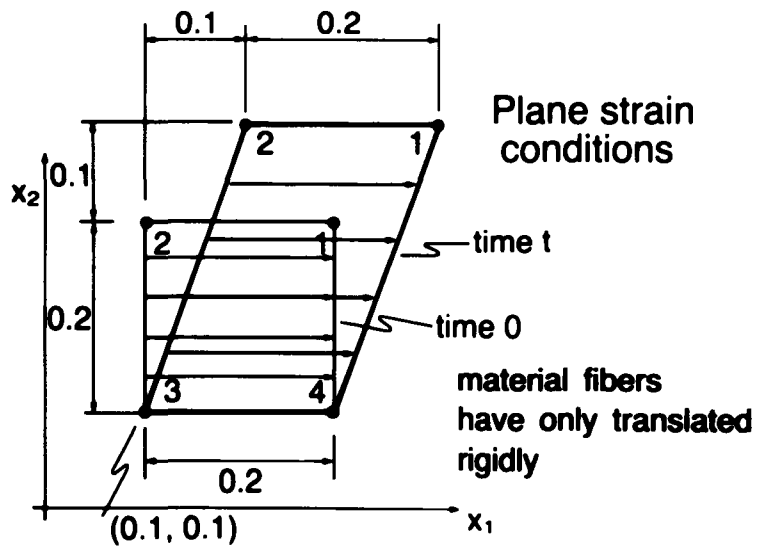
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Example: Calculation of  ${}^0\underline{B}_L, {}^0\underline{B}_{NL}$

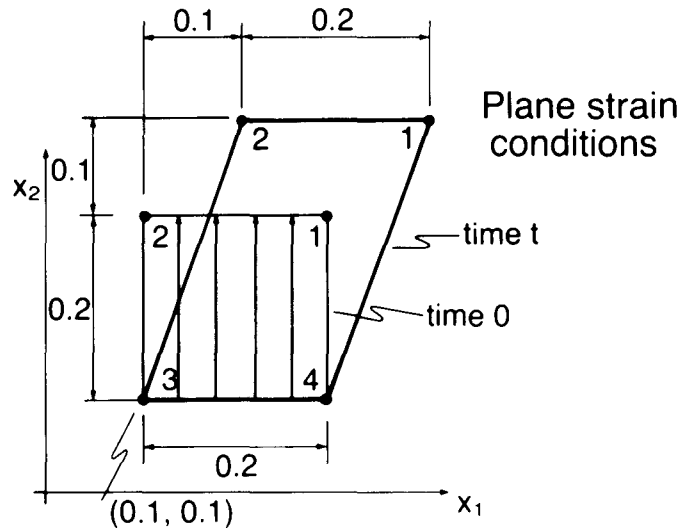


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Example: Calculation of  ${}^0\underline{B}_L, {}^0\underline{B}_{NL}$

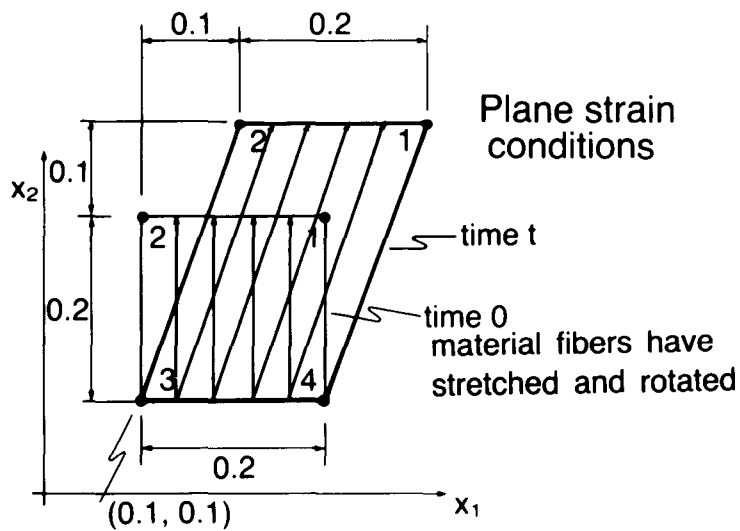


Example: Calculation of  ${}^t_0\underline{B}_L, {}^t_0\underline{B}_{NL}$



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Example: Calculation of  ${}^t_0\underline{B}_L, {}^t_0\underline{B}_{NL}$

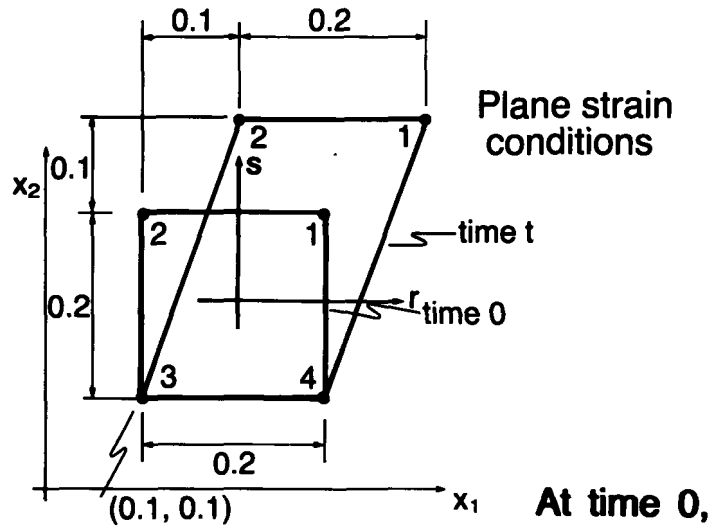


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Example: Calculation of  ${}^0\underline{B}_L$ ,  ${}^t\underline{B}_{NL}$



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We can now perform a Jacobian transformation between the  $(r, s)$  coordinate system and the  $({}^0x_1, {}^0x_2)$  coordinate system:

$$\text{By inspection, } \frac{\partial {}^0x_1}{\partial r} = 0.1, \frac{\partial {}^0x_2}{\partial r} = 0$$

$$\frac{\partial {}^0x_1}{\partial s} = 0, \frac{\partial {}^0x_2}{\partial s} = 0.1$$

$$\text{Hence } {}^0\underline{J} = \begin{bmatrix} 0.1 & 0 \\ 0 & 0.1 \end{bmatrix}, |{}^0\underline{J}| = 0.01$$

$$\text{and } \frac{\partial}{\partial {}^0x_1} = 10 \frac{\partial}{\partial r}, \frac{\partial}{\partial {}^0x_2} = 10 \frac{\partial}{\partial s}$$

Now we use the interpolation functions to compute  ${}^t_0u_{1,1}$ ,  ${}^t_0u_{1,2}$ :

node k	$\frac{\partial h_k}{\partial^0x_1}$	$\frac{\partial h_k}{\partial^0x_2}$	${}^t u_1^k$	$\frac{\partial h_k}{\partial^0x_1} {}^t u_1^k$	$\frac{\partial h_k}{\partial^0x_2} {}^t u_1^k$
1	$2.5(1 + s)$	$2.5(1 + r)$	0.1	$0.25(1 + s)$	$0.25(1 + r)$
2	$-2.5(1 + s)$	$2.5(1 - r)$	0.1	$-0.25(1 + s)$	$0.25(1 - r)$
3	$-2.5(1 - s)$	$-2.5(1 - r)$	0.0	0	0
4	$2.5(1 - s)$	$-2.5(1 + r)$	0.0	0	0

$$\text{Sum: } \underbrace{0.0}_{{}^t_0u_{1,1}} \quad \underbrace{0.5}_{{}^t_0u_{1,2}}$$

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For this simple problem, we can compute the displacement derivatives by inspection:

From the given dimensions,

$${}^t_0\underline{X} = \begin{bmatrix} 1.0 & 0.5 \\ 0.0 & 1.5 \end{bmatrix}$$

Hence

$$\begin{aligned} {}^t_0u_{1,1} &= {}^tX_{11} - 1 = 0 \\ {}^t_0u_{1,2} &= {}^tX_{12} = 0.5 \\ {}^t_0u_{2,1} &= {}^tX_{21} = 0 \\ {}^t_0u_{2,2} &= {}^tX_{22} - 1 = 0.5 \end{aligned}$$

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We can now construct the columns in  ${}^t\underline{B}_L$  that correspond to node 3:

$$\left[ \begin{array}{c|c|c|c} \dots & -2.5(1-s) & 0 & \dots \\ & 0 & -2.5(1-r) & \\ & -2.5(1-r) & -2.5(1-s) & \end{array} \right] {}^t\underline{B}_{L0}$$

$$\left[ \begin{array}{c|c|c|c} \dots & 0 & 0 & \dots \\ & -1.25(1-r) & -1.25(1-r) & \\ & -1.25(1-s) & -1.25(1-s) & \end{array} \right] {}^t\underline{B}_{L1}$$

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Similarly, we construct the columns in  ${}^t\underline{B}_{NL}$  that correspond to node 3:

$$\left[ \begin{array}{c|c|c|c} \dots & -2.5(1-s) & 0 & \dots \\ & -2.5(1-r) & 0 & \\ & 0 & -2.5(1-s) & \\ & 0 & -2.5(1-r) & \end{array} \right]$$

Consider next the element matrices required for the updated Lagrangian formulation:

Element Matrix	Matrices Required
$\underline{{}^t\mathbf{K}}_L$	$\underline{{}^t\mathbf{C}}$ , $\underline{{}^t\mathbf{B}}_L$
$\underline{{}^t\mathbf{K}}_{NL}$	$\underline{{}^t\mathbf{T}}$ , $\underline{{}^t\mathbf{B}}_{NL}$
$\underline{{}^t\mathbf{F}}$	$\underline{{}^t\hat{\mathbf{T}}}$ , $\underline{{}^t\mathbf{B}}_L$

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We define  $\underline{{}^t\mathbf{C}}$  so that

$$\begin{bmatrix} {}^t\mathbf{S}_{11} \\ {}^t\mathbf{S}_{22} \\ {}^t\mathbf{S}_{12} \\ {}^t\mathbf{S}_{33} \end{bmatrix} = \underline{{}^t\mathbf{C}} \begin{bmatrix} {}^t\mathbf{e}_{11} \\ {}^t\mathbf{e}_{22} \\ 2\,{}^t\mathbf{e}_{12} \\ {}^t\mathbf{e}_{33} \end{bmatrix} \quad \text{analogous to} \quad {}^t\mathbf{S}_{ij} = {}^t\mathbf{C}_{ij,rs} \, {}^t\mathbf{e}_{rs}$$

For example, we may choose (axisymmetric analysis),

$$\underline{{}^t\mathbf{C}} = \frac{E(1-\nu)}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1 & \frac{\nu}{1-\nu} & 0 & \frac{\nu}{1-\nu} \\ \frac{\nu}{1-\nu} & 1 & 0 & \frac{\nu}{1-\nu} \\ 0 & 0 & \frac{1-2\nu}{2(1-\nu)} & 0 \\ \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & 0 & 1 \end{bmatrix}$$

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**Transparency  
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We note that the incremental strain components are, in two-dimensional analysis,

$${}^t\epsilon_{11} = \frac{\partial u_1}{\partial {}^tX_1} = {}^t u_{1,1}$$

$${}^t\epsilon_{22} = {}^t u_{2,2}$$

$$2 {}^t\epsilon_{12} = {}^t u_{1,2} + {}^t u_{2,1}$$

$${}^t\epsilon_{33} = u_1 / {}^tX_1$$

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and

$${}^t\eta_{11} = \frac{1}{2} (({}^t u_{1,1})^2 + ({}^t u_{2,1})^2)$$

$${}^t\eta_{22} = \frac{1}{2} (({}^t u_{1,2})^2 + ({}^t u_{2,2})^2)$$

$${}^t\eta_{12} = {}^t\eta_{21} = \frac{1}{2} ({}^t u_{1,1} {}^t u_{1,2} + {}^t u_{2,1} {}^t u_{2,2})$$

$${}^t\eta_{33} = \frac{1}{2} \left( \frac{u_1}{{}^tX_1} \right)^2$$

We construct  $\underline{tB}_L$  so that

$$\begin{bmatrix} \underline{t}e_{11} \\ \underline{t}e_{22} \\ 2 \underline{t}e_{12} \\ \underline{t}e_{33} \end{bmatrix} = \underline{t}e = \underline{tB}_L \hat{u}$$

only included for axisymmetric analysis

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Entries in  $\underline{tB}_L$ :

		node k			
		$u_1^k$	$u_2^k$		
...	$\underline{t}h_{k,1}$	0	0	$\underline{t}h_{k,2}$	...
...	0	$\underline{t}h_{k,2}$	$\underline{t}h_{k,1}$	$\underline{t}h_{k,1}$	...
	$h_k/\underline{t}x_1$	0			

$$\begin{bmatrix} \vdots \\ \underline{u}_1^k \\ \underline{u}_2^k \\ \vdots \end{bmatrix}$$

↑  
node k  
↓

only included for axisymmetric analysis

This is similar in form to the  $\underline{B}$  matrix used in linear analysis.

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We construct  ${}^t\mathbf{B}_{NL}$  and  ${}^t\mathbf{T}$  so that

$$\delta \hat{\mathbf{u}}^T {}^t\mathbf{B}_{NL}^T {}^t\mathbf{T} {}^t\mathbf{B}_{NL} \hat{\mathbf{u}} = {}^t\mathbf{T}_{ij} \delta \epsilon_{ij}$$

Entries in  ${}^t\mathbf{T}$ :

${}^tT_{11}$	${}^tT_{12}$	0	0	0	included only for axisymmetric analysis
${}^tT_{21}$	${}^tT_{22}$	0	0	0	
0	0	${}^tT_{11}$	${}^tT_{12}$	0	
0	0	${}^tT_{21}$	${}^tT_{22}$	0	
0	0	0	0	${}^tT_{33}$	

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Entries in  ${}^t\mathbf{B}_{NL}$ :

		node k			
		$u_1^k$	$u_2^k$		
...	${}^th_{k,1}$	0	...	$\vdots$	node k
...	${}^th_{k,2}$	0	...	$u_1^k$	
...	0	${}^th_{k,1}$	...	$u_2^k$	
...	0	${}^th_{k,2}$	...		
		$h_k/x_1$	0		

included only for  
axisymmetric analysis

${}^t\hat{\underline{T}}$  is constructed so that

$$\delta \underline{\hat{u}}^T {}^t\underline{B}_L^T {}^t\hat{\underline{T}} = {}^t\underline{\tau}_{ij} \delta {}^t e_{ij}$$

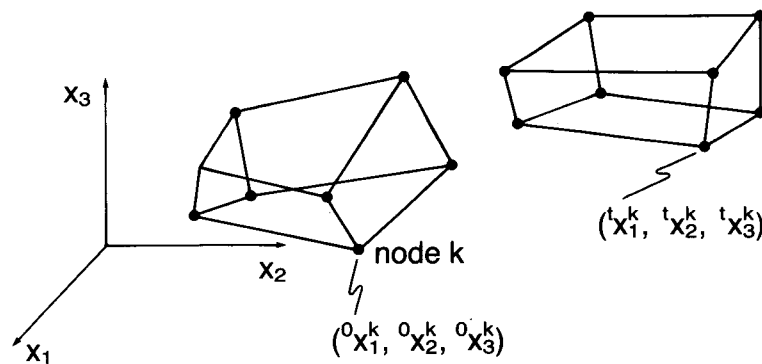
Entries in  ${}^t\hat{\underline{T}}$ :

$$\begin{bmatrix} {}^tT_{11} \\ {}^tT_{22} \\ {}^tT_{12} \\ {}^tT_{33} \end{bmatrix}$$

included only for axisymmetric analysis

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Three-dimensional elements



Transparency 7-44



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Here we now use

$${}^0x_1 = \sum_{k=1}^N h_k {}^0x_1^k, \quad {}^0x_2 = \sum_{k=1}^N h_k {}^0x_2^k$$
$${}^0x_3 = \sum_{k=1}^N h_k {}^0x_3^k,$$

where the  $h_k$ 's are the isoparametric interpolation functions of the three-dimensional element.

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Also

$${}^tx_1 = \sum_{k=1}^N h_k {}^tx_1^k, \quad {}^tx_2 = \sum_{k=1}^N h_k {}^tx_2^k$$
$${}^tx_3 = \sum_{k=1}^N h_k {}^tx_3^k$$

and then all the concepts and derivations already discussed are directly applicable to the derivation of the three-dimensional element matrices.

MIT OpenCourseWare  
<http://ocw.mit.edu>

**Resource: Finite Element Procedures for Solids and Structures**  
Klaus-Jürgen Bathe

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