

PROFESSOR: Let's look at the magnitude squared of those waves that we've already defined here. We have two solutions, one for no potential and one for a real potential. Both are for some finite range potential.

We have $\phi(x)^2$ is equal to $\sin^2(kx)$. And $\psi(x)^2$ is equal to $\sin^2(kx + \delta)$. Even from this information you can get something.

Think of x plotted here. Here's $x = 0$. And there's this \sin^2 . This wave for $\phi(x)^2$.

Suppose you're looking at some feature-- a maximum, a minimum-- of this function. Suppose the feature happens when the argument, kx , is equal to some number, a_0 . Whatever feature-- this number a_0 could be 0, in which case you're looking at a minimum, it could be $\pi/2$, in which case you're looking at a maximum-- some feature of \sin^2 .

Well the same feature will appear in this case when the whole argument is equal to a_0 . So while this one happens at $x = a_0/k$, here it happens at $x = a_0/k - \delta/k$. If this is the probability density associated to the solution for no potential and it has a maximum here, the maximum of the true solution-- say, here-- would appear at a distance equal to δ/k . Earlier-- so this is like the x , and this is like the \tilde{x} -- that feature would appear, δ/k in that direction.

So this is ψ . This is ψ^2 . So we conclude, for example, that when δ is greater than 0, the wave is pulled. $\delta = 0$, the two shapes are on top of each other. For δ different from 0, the wave function is pulled in. So $\delta > 0$, ψ is pulled in.

What could we think of this? The potential is attractive. It's pulling in the wave function. Attractive.

$\delta < 0$, the wave is pushed out. It would be in the other direction, and the ψ is pushed out. Potential is repulsive. So a little bit of information even from the signs of this thing.

We want to define one last thing, and then we'll stop. It's the concept of the scattered wave. What should we call the scattered wave? We will define the scattered wave ψ_s as the extra piece in the solution-- the ψ solution-- that would vanish without potential.

So we say, you have a ψ , but if you didn't have a potential, what part of this ψ would survive? Think of writing the ψ of x as the solution without the potential plus the extra part, the scattered wave, ψ_s . So this is the definition.

The full scattering solution, the full solution when you have a potential, can be written in a solution without the potential and this scattered thing. Now, you may remember-- we just did it a second ago-- that this original solution and the ψ solution have the same incoming wave. The incoming wave up there is the same for the ψ solution as for this one.

So the incoming waves are the same. So only the outgoing waves are different. And this represents how much more of an outgoing wave you get than from what you would have gotten with ψ . So this must be an outgoing wave.

We'll just plug in the formula here. ψ_s is equal to ψ minus ϕ . And it's equal to $\frac{1}{2i} e^{ikx} + 2i\delta \sin \delta - e^{-ikx} - \frac{1}{2i}$ -- the ϕ -- into the ikx minus e^{-ikx} . So the incoming waves were the same. Indeed, they cancelled. But the outgoing waves are not. You can factor an e^{ikx} , and you get $e^{2i\delta} \sin \delta - 1$. Which is equal to $e^{ikx} e^{i\delta} \sin \delta$.

There we go. We have the answer for the scattered wave. It's proportional to $\sin \delta$, which, again, makes sense. If δ is equal to 0, there is no scattering. It's an outgoing wave and all is good. So I'll write it like this. ψ_s is equal to $A e^{ikx}$, with A equal to $e^{i\delta} \sin \delta$.