Class 30: Outline

Hour 1:

Traveling & Standing Waves

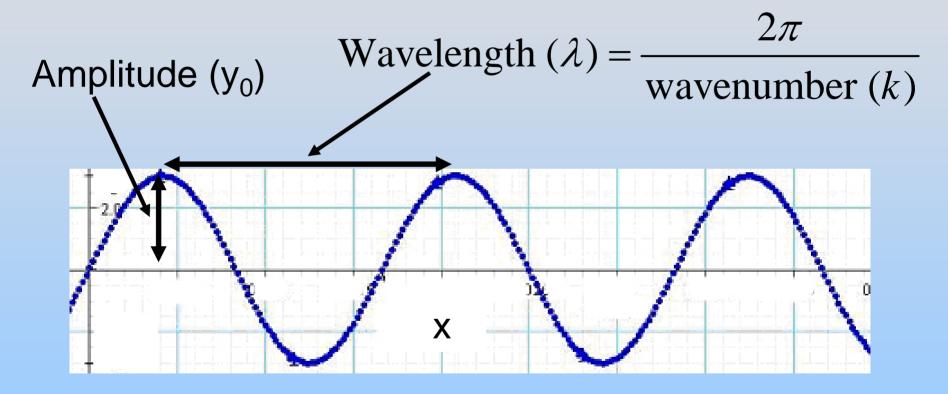
Hour 2:

Electromagnetic (EM) Waves

Last Time: Traveling Waves

Traveling Sine Wave

Now consider $f(x) = y = y_0 \sin(kx)$:

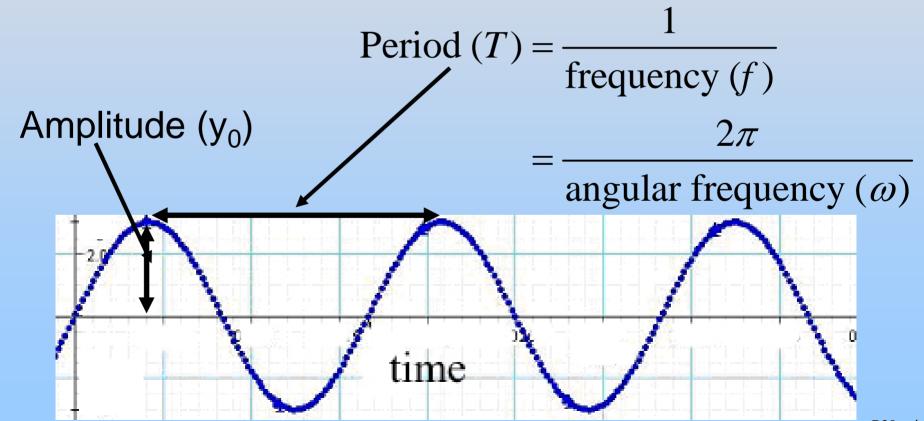


What is g(x,t) = f(x+vt)? Travels to left at velocity $v = y_0 \sin(k(x+vt)) = y_0 \sin(kx+kvt)$

Traveling Sine Wave

$$y = y_0 \sin(kx + kvt)$$

At x=0, just a function of time: $y = y_0 \sin(kvt) \equiv y_0 \sin(\omega t)$



Traveling Sine Wave

- Wavelength: λ
- Frequency : f

$$y = y_0 \sin(kx - \omega t)$$

- Wave Number: $k = \frac{2\pi}{\lambda}$
- Angular Frequency: $\omega = 2\pi f$
- Period: $T = \frac{1}{f} = \frac{2\pi}{\omega}$
- Speed of Propagation: $v = \frac{\omega}{k} = \lambda f$
- Direction of Propagation: +x

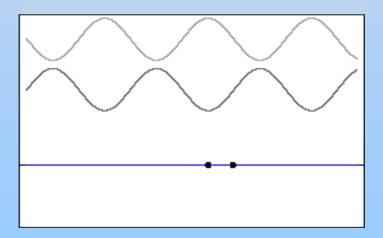
This Time: Standing Waves

Standing Waves

What happens if two waves headed in opposite directions are allowed to interfere?

$$E_1 = E_0 \sin(kx - \omega t) \qquad E_2 = E_0 \sin(kx + \omega t)$$

Superposition:
$$E = E_1 + E_2 = 2E_0 \sin(kx) \cos(\omega t)$$

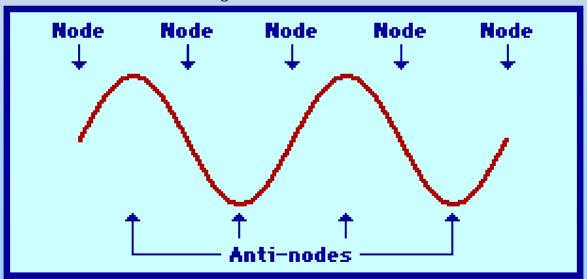


Standing Waves: Who Cares?

Most commonly seen in resonating systems:

Musical Instruments, Microwave Ovens

$$E = 2E_0 \sin(kx) \cos(\omega t)$$



Standing Waves: Bridge

Tacoma Narrows Bridge Oscillation:

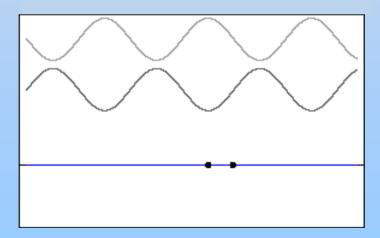
http://www.pbs.org/wgbh/nova/bridge/tacoma3.html

Group Work: Standing Waves

Do Problem 2

$$E_1 = E_0 \sin(kx - \omega t) \qquad E_2 = E_0 \sin(kx + \omega t)$$

Superposition: $E = E_1 + E_2 = 2E_0 \sin(kx) \cos(\omega t)$



Last Time: Maxwell's Equations

Maxwell's Equations

$$\iint_{S} \vec{\mathbf{E}} \cdot d\vec{\mathbf{A}} = \frac{Q_{in}}{\mathcal{E}_{0}}$$

(Gauss's Law)

$$\oint_C \vec{\mathbf{E}} \cdot d \vec{\mathbf{s}} = -\frac{d\Phi_B}{dt}$$

(Faraday's Law)

$$\iint_{S} \vec{\mathbf{B}} \cdot d\vec{\mathbf{A}} = 0$$

(Magnetic Gauss's Law)

$$\oint_C \vec{\mathbf{B}} \cdot d \vec{\mathbf{s}} = \mu_0 I_{enc} + \mu_0 \varepsilon_0 \frac{d\Phi_E}{dt}$$

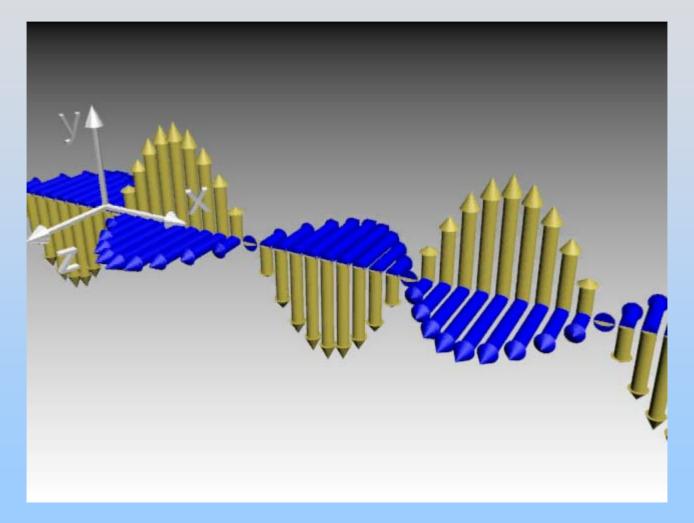
(Ampere-Maxwell Law)

$$\vec{\mathbf{F}} = q(\vec{\mathbf{E}} + \vec{\mathbf{v}} \times \vec{\mathbf{B}})$$

(Lorentz force Law)

Which Leads To... EM Waves

Electromagnetic Radiation: Plane Waves



Traveling E & B Waves

- Wavelength: λ
- Frequency : f

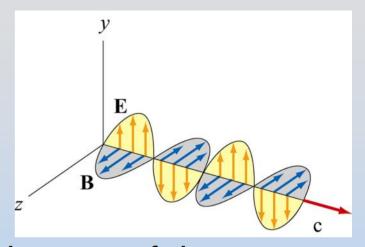
$$\vec{\mathbf{E}} = \hat{\mathbf{E}}E_0 \sin(kx - \omega t)$$

- Wave Number: $k = \frac{2\pi}{\lambda}$
- Angular Frequency: $\omega = 2\pi f$
- Period: $T = \frac{1}{f} = \frac{2\pi}{\omega}$
- Speed of Propagation: $v = \frac{\omega}{k} = \lambda f$
- Direction of Propagation: +x

Properties of EM Waves

Travel (through vacuum) with

speed of light
$$v = c = \frac{1}{\sqrt{\mu_0 \varepsilon_0}} = 3 \times 10^8 \frac{m}{s}$$



At every point in the wave and any instant of time, E and B are in phase with one another, with

$$\frac{E}{B} = \frac{E_0}{B_0} = c$$

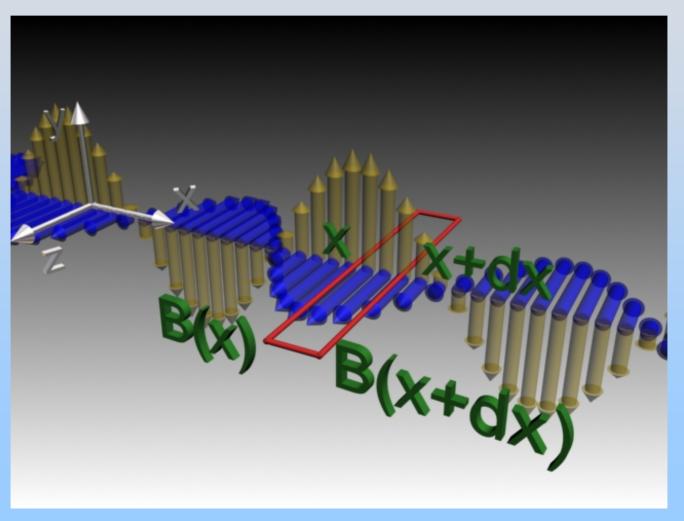
E and B fields perpendicular to one another, and to the direction of propagation (they are transverse):

Direction of propagation = Direction of $\vec{E} \times \vec{B}$

PRS Questions: Direction of Propagation

How Do Maxwell's Equations Lead to EM Waves? Derive Wave Equation

Start with Ampere-Maxwell Eq: $\oint_C \vec{\mathbf{B}} \cdot d\vec{\mathbf{s}} = \mu_0 \varepsilon_0 \frac{d}{dt} \int \vec{\mathbf{E}} \cdot d\vec{\mathbf{A}}$

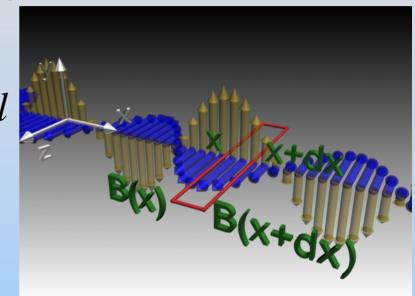


Start with Ampere-Maxwell Eq: $\oint_{C} \vec{\mathbf{B}} \cdot d\vec{\mathbf{s}} = \mu_0 \varepsilon_0 \frac{d}{dt} \int_{C} \vec{\mathbf{E}} \cdot d\vec{\mathbf{A}}$

Apply it to red rectangle:

$$\oint_C \vec{\mathbf{B}} \cdot d\vec{\mathbf{s}} = B_z(x,t)l - B_z(x+dx,t)l$$

$$\mu_0 \varepsilon_0 \frac{d}{dt} \int \vec{\mathbf{E}} \cdot d\vec{\mathbf{A}} = \mu_0 \varepsilon_0 \left(l \, dx \frac{\partial E_y}{\partial t} \right)$$



$$-\frac{B_z(x+dx,t)-B_z(x,t)}{dx} = \mu_0 \varepsilon_0 \frac{\partial E_y}{\partial t}$$

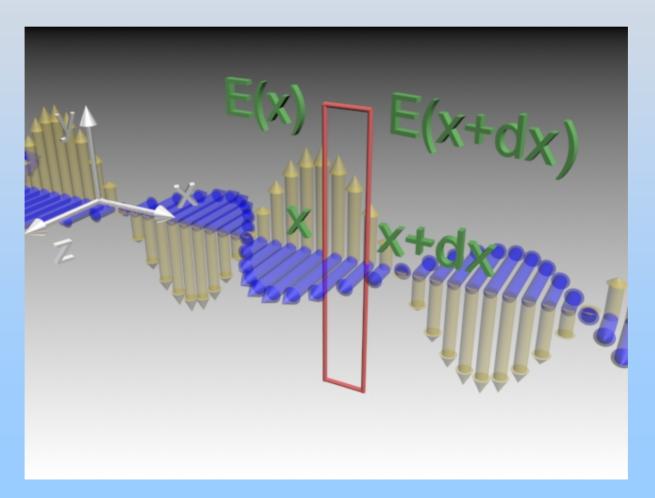
So in the limit that dx is very small:

$$-\frac{\partial B_z}{\partial x} = \mu_0 \varepsilon_0 \frac{\partial E_y}{\partial t}$$

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Now go to Faraday's Law

$$\oint_C \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}} = -\frac{d}{dt} \int_C \vec{\mathbf{B}} \cdot d\vec{\mathbf{A}}$$



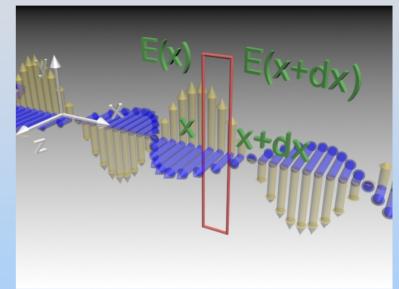
Faraday's Law:

$$\oint_C \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}} = -\frac{d}{dt} \int_C \vec{\mathbf{B}} \cdot d\vec{\mathbf{A}}$$

Apply it to red rectangle:

$$\oint_C \vec{\mathbf{E}} \cdot d \vec{\mathbf{s}} = E_y(x + dx, t)l - E_y(x, t)l$$

$$-\frac{d}{dt}\int \vec{\mathbf{B}} \cdot d\vec{\mathbf{A}} = -ldx \frac{\partial B_z}{\partial t}$$



$$\frac{E_{y}(x+dx,t) - E_{y}(x,t)}{dx} = -\frac{\partial B_{z}}{\partial t}$$

So in the limit that dx is very small:

$$\frac{\partial E_{y}}{\partial x} = -\frac{\partial B_{z}}{\partial t}$$

1D Wave Equation for E

$$\frac{\partial E_{y}}{\partial x} = -\frac{\partial B_{z}}{\partial t} \qquad -\frac{\partial B_{z}}{\partial x} = \mu_{0} \varepsilon_{0} \frac{\partial E_{y}}{\partial t}$$

Take x-derivative of 1st and use the 2nd equation

$$\frac{\partial}{\partial x} \left(\frac{\partial E_{y}}{\partial x} \right) = \frac{\partial^{2} E_{y}}{\partial x^{2}} = \frac{\partial}{\partial x} \left(-\frac{\partial B_{z}}{\partial t} \right) = -\frac{\partial}{\partial t} \left(\frac{\partial B_{z}}{\partial x} \right) = \mu_{0} \varepsilon_{0} \frac{\partial^{2} E_{y}}{\partial t^{2}}$$

$$\left| \frac{\partial^2 E_y}{\partial x^2} = \mu_0 \varepsilon_0 \frac{\partial^2 E_y}{\partial t^2} \right|$$

1D Wave Equation for E

$$\frac{\partial^2 E_y}{\partial x^2} = \mu_0 \varepsilon_0 \frac{\partial^2 E_y}{\partial t^2}$$

This is an equation for a wave. Let: $E_y = f(x - vt)$

$$\frac{\partial^{2} E_{y}}{\partial x^{2}} = f''(x - vt)$$

$$\frac{\partial^{2} E_{y}}{\partial t^{2}} = v^{2} f''(x - vt)$$

$$v^{2} = \frac{1}{\mu_{0} \mathcal{E}_{0}}$$

1D Wave Equation for B

$$\frac{\partial B_{z}}{\partial t} = -\frac{\partial E_{y}}{\partial x} \qquad \frac{\partial B_{z}}{\partial x} = -\mu_{0} \varepsilon_{0} \frac{\partial E_{y}}{\partial t}$$

Take x-derivative of 1st and use the 2nd equation

$$\frac{\partial}{\partial t} \left(\frac{\partial B_z}{\partial t} \right) = \frac{\partial^2 B_z}{\underline{\partial t^2}} = \frac{\partial}{\partial t} \left(-\frac{\partial E_y}{\partial x} \right) = -\frac{\partial}{\partial x} \left(\frac{\partial E_y}{\partial t} \right) = \frac{1}{\underline{\mu_0 \varepsilon_0}} \frac{\partial^2 B_z}{\partial x^2}$$

$$\frac{\partial^2 B_z}{\partial x^2} = \mu_0 \varepsilon_0 \frac{\partial^2 B_z}{\partial t^2}$$

Electromagnetic Radiation

Both E & B travel like waves:

$$\frac{\partial^{2} E_{y}}{\partial x^{2}} = \mu_{0} \varepsilon_{0} \frac{\partial^{2} E_{y}}{\partial t^{2}} \qquad \frac{\partial^{2} B_{z}}{\partial x^{2}} = \mu_{0} \varepsilon_{0} \frac{\partial^{2} B_{z}}{\partial t^{2}}$$

But there are strict relations between them:

$$\frac{\partial B_{z}}{\partial t} = -\frac{\partial E_{y}}{\partial x} \qquad \qquad \frac{\partial B_{z}}{\partial x} = -\mu_{0} \varepsilon_{0} \frac{\partial E_{y}}{\partial t}$$

Here, E_y and B_z are "the same," traveling along x axis

Amplitudes of E & B

Let
$$E_v = E_0 f(x - vt)$$
; $B_z = B_0 f(x - vt)$

$$\frac{\partial B_z}{\partial t} = -\frac{\partial E_y}{\partial x} \Longrightarrow -vB_0 f'(x - vt) = -E_0 f'(x - vt)$$

$$\Rightarrow vB_0 = E_0$$

E_v and B_z are "the same," just different amplitudes

Group Problem: EM Standing Waves

Consider EM Wave approaching a perfect conductor:

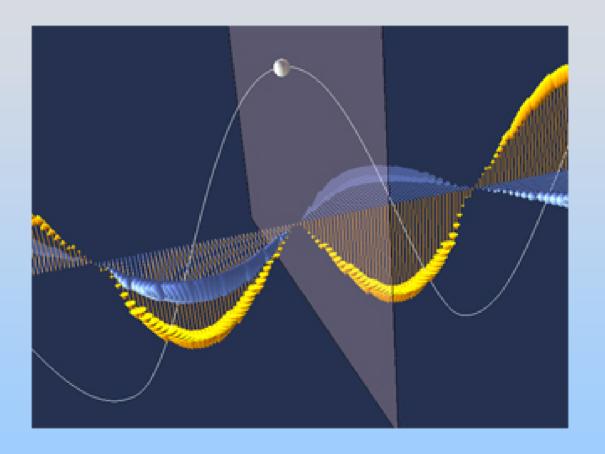
$$\vec{\mathbf{E}}_{\text{incident}} = \hat{x}E_0\cos(kz - \omega t)$$

 $\mathbf{E}_{\text{incident}} = \hat{x}E_0\cos(kz - \omega t)$ If the conductor fills the XY plane at Z=0 then the wave will reflect and add to the incident wave

- 1. What must the total E field ($E_{inc}+E_{ref}$) at Z=0 be?
- 2. What is $E_{reflected}$ for this to be the case?
- 3. What are the accompanying B fields? (B_{inc} & B_{ref})
- 4. What are E_{total} and B_{total} ? What is B(Z=0)?
- 5. What current must exist at Z=0 to reflect the wave? Give magnitude and direction.

Recall:
$$\cos(A+B) = \cos(A)\cos(B) - \sin(A)\sin(B)$$

Next Time: How Do We Generate Plane Waves?



http://ocw.mit.edu/ans7870/8/8.02T/f04/visualizations/light/09-planewaveapp/09-planewaveapp320.html