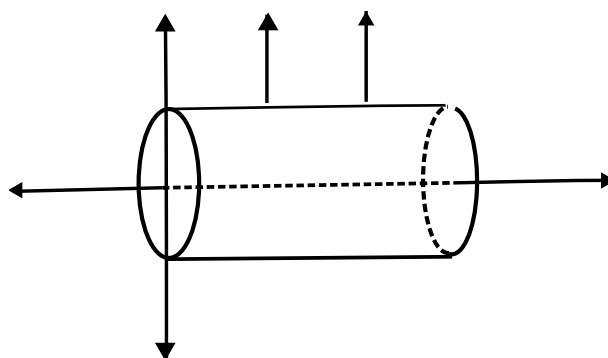


8.022 Lecture Notes Class 9 - 09/20/2006

Given line of charge with density  $\lambda$

Line Charge

Find  $\vec{E}(\vec{r})$  everywhere



$$\int_S \vec{E} \cdot d\vec{a} = \frac{1}{\epsilon_0} Q_{enclosed} \quad Q_{enclosed} = \lambda L$$

$$L 2\pi r \cdot \vec{E}(\vec{r}) = \frac{\lambda L}{\epsilon_0}$$

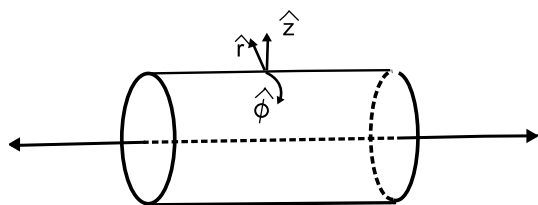
$$\vec{E}(\vec{r}) = \frac{\lambda}{2\pi\epsilon_0 r} \hat{r}$$

No  $\hat{\phi}$  because  $\vec{\nabla} \times \vec{E} = 0$ .

No  $\hat{z}$  because vectors cancel out

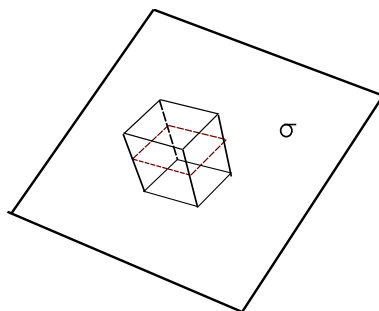
$$E(r) \propto \frac{1}{r}$$

Like a plane : Cross sections reveal that field lines spread only perpendicular to line.



$$E(r) \propto \frac{1}{r}$$

N-dimensional  $\rightarrow \frac{1}{r^{N-1}}$  (Surface area of N-D object is (N-1)D)



## Plane Charge

$$Q_{enclosed} = \sigma \cdot A^2$$

$$\int \vec{E} \cdot d\vec{A} = \frac{Q_{enclosed}}{\epsilon_0}$$

$$\iiint_S \vec{E} \cdot dA = \frac{\sigma \cdot A^2}{\epsilon_0}$$

$$\int_0^A \int_0^A \hat{z} \vec{E}(\vec{r}) dx dy + \int_0^A \int_0^A -\vec{E}(\vec{r})(-\hat{z}) dx dy = \frac{\sigma \cdot A^2}{\epsilon_0}$$

$$2A^2 E_z(\vec{r}) = \frac{\sigma A^2}{\epsilon_0}$$

$$E_z(\vec{r}) = \frac{\sigma}{2\epsilon_0} \vec{E} = \hat{n} \cdot \frac{\sigma}{2\epsilon_0}$$

## Common Electric Fields

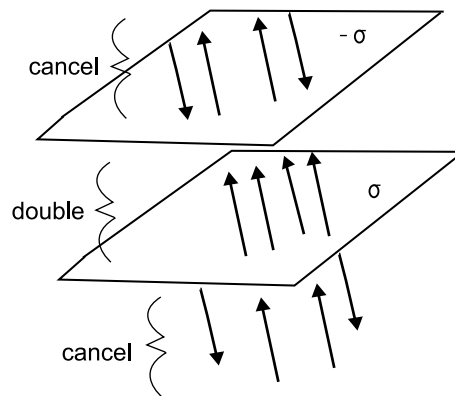
Line:

$$\vec{E}(\vec{r}) = \frac{\lambda}{2\pi\epsilon_0 r} \hat{r}$$

Plane:

$$\vec{E}(\vec{r}) = \hat{n} \cdot \frac{\sigma}{2\epsilon_0}$$

Double Plane:



$$\vec{E}(\vec{r}) = \hat{n} \cdot \frac{\sigma}{\epsilon} \quad (\text{inside})$$

$$\vec{E}(\vec{r}) = \vec{0} \quad (\text{outside})$$

$$\int_{\vec{a}}^{\vec{b}} \vec{E} \cdot d\vec{l} \quad \vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

$$= \frac{1}{4\pi\epsilon_0} \int_{\vec{a}}^{\vec{b}} \frac{q}{r^2} dr$$

$$= -\frac{1}{4\pi\epsilon_0} \left( \frac{q}{r(\vec{b})} - \frac{q}{r(\vec{a})} \right)$$

$$\oint \vec{E} \cdot d\vec{l} = 0 \quad \text{because } r(\vec{b}) = r(\vec{a})$$

$$\Rightarrow \vec{\nabla} \times \vec{E} = \vec{0}$$