#### 22.615, MHD Theory of Fusion Systems Prof. Freidberg Lecture 11: Flux Conserving Tokamak - Con'd

## A Simple Approximation

- 1. Instead of choosing  $F(\psi)$  so  $q(\psi)$  is the same everywhere, we choose a simpler  $F(\psi)$  so that only q(0) and q(a) remain the same (as  $\beta_t$  increases).
- 2. Choose  $dp/d\psi = \text{const}$ ,  $dF^2/d\psi = \text{const}$ . This is the same model we have already investigated.
- 3. The model has only two free parameters:  $A, C \rightarrow \beta_t, q_*$ .
- 4. Thus, as  $\beta_t$  increases, there is only one degree of freedom,  $q_*$ , remaining.
- 5. Therefore, we cannot adjust  $q_*$  so that both  $q_0$  and  $q_a$  remain fixed: this would be an overdetermined system.
- 6. We make an ultra simple approximation and choose  $q_*$  so that only  $q_a$  remains fixed. This prevents the formation of a separatrix which requires  $q_a \rightarrow \infty$ .

## HBT Equilibrium

$$\mu_{0}\rho = \beta_{t}B_{0}^{2}\left(1-\rho^{2}\right)\left[1-v\rho\cos\theta\right]$$

$$B_{\theta} = \frac{\epsilon}{q_{\star}}\left[\rho + \frac{v}{2}\left(3\rho^{2}-1\right)\cos\theta\right]$$

$$\hat{B}_{\theta} = \frac{\epsilon}{q_{\star}}\left[\frac{1}{\rho} + \frac{v}{2}\left(1+\frac{1}{\rho^{2}}\right)\cos\theta\right]$$

$$q_{a} = \frac{q_{\star}}{\left(1-v^{2}\right)^{1/2}}$$

$$v = \frac{\beta_{t}q_{\star}^{2}}{\epsilon}$$

$$\rho = r/a$$

- 1. HBT: Express all quantities in terms of  $\beta_t$ ,  $q_* \sim 1/I$
- 2. FCT: Express all quantities in terms of  $\beta_t$ ,  $q_a$  (held fixed). Examine the behavior as  $\beta_t$  increases. Are there any equilibrium limits?

### Procedure

- 1. Define  $v_* = \beta_t q_a^2 / \epsilon \propto \beta_t$  since  $q_a$  is held fixed in the FCT.
- 2.  $v_*$  is the heating parameter: as  $\beta_t$  increases,  $v_*$  increases.
- 3. For the HBT:  $\nu = \beta_t q_*^2 / \epsilon \propto \beta_t$  for fixed *I*.
- 4. v is the heating parameter for fixed *I*: as  $\beta_t$  increases, v increases.

Relation between v and  $v_*$ 

1. 
$$v = \frac{\beta_t q_\star^2}{\epsilon} = \frac{\beta_t q_a^2}{\epsilon} \frac{q_\star^2}{q_a^2} = v_\star (1 - v^2)$$
  
2.  $v^2 + \frac{v}{v_\star} - 1 = 0$   
 $v = \frac{2v_\star}{(1 + 4v_\star^2)^{1/2} + 1}$ 

#### Compute the Physical Quantities in Terms of $v_*$ and Compare with the HBT

1. 
$$I \propto 1/q_{\star}$$
  
a. HBT:  $\frac{1}{q_{\star}} = \text{const.}$  fixed *I*  
b. FCT:  $\frac{1}{q_{\star}} = \frac{1}{q_{a}} \frac{1}{(1-v^{2})^{1/2}} = \frac{1}{q_{a}} \left(\frac{v_{\star}}{v}\right)^{1/2}$   
 $\frac{1}{q_{\star}} = \frac{1}{q_{a}} \left[\frac{1+(1+4v_{\star}^{2})^{1/2}}{2}\right]^{1/2}$   
2.  $B_{v}$   
a. HBT:  $B_{v} = \frac{\mu_{0}I}{4\pi R_{0}} \beta_{p} = \frac{\epsilon B_{0}}{q_{\star}} \frac{\epsilon \beta_{p}}{2} = \frac{\epsilon B_{0}}{2} \frac{v}{q_{\star}}$   
b. FCT:  $B_{v} = \frac{\epsilon B_{0}}{2} \frac{v}{q_{\star}} = \frac{\epsilon B_{0}}{2} \frac{1}{q_{a}} \left[\frac{1+(1+4v_{\star}^{2})^{1/2}}{2}\right]^{1/2} \frac{2v_{\star}}{1+(1+4v_{\star}^{2})^{1/2}}$ 

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$$B_{\nu} = \frac{\epsilon B_0}{2} \frac{\nu_*}{q_a} \left[ \frac{2}{1 + (1 + 4\nu_*^2)^{1/2}} \right]^{1/2}$$

3. *ρ*<sub>s</sub>

a. HBT: 
$$\rho_s = \frac{1}{\nu} \left[ 1 + \left( 1 - \nu^2 \right)^{1/2} \right]$$
  
b. FCT:  $\rho_s = \frac{1 + \left( 1 + 4\nu_*^2 \right)^{1/2}}{2\nu_*} \left[ 1 + \left( \frac{2}{1 + \left( 1 + 4\nu_*^2 \right)^{1/2}} \right)^{1/2} \right]$ 

4. Define the plasma evolution in  $\beta_t - q_{\star}$  space as  $\beta_t$  increases

a. HBT: 
$$\frac{\beta_t q_*^2}{\epsilon} = v$$
  
 $q_* = \text{const.}$   
b. FCT:  $\frac{\beta_t q_*^2}{\epsilon} = v_*$  (1)  
 $\frac{1}{q_*} = \frac{1}{q_a} \left[ \frac{1 + (1 + 4v_*^2)^{1/2}}{2} \right]^{1/2}$  (2)

c. Solve (2) for  $v_*$  and substitute into (1) to give  $\beta_t = F(q_*)$ 

$$v_{\star}^{2} = \frac{q_{a}^{2}}{q_{\star}^{2}} \left[ \frac{q_{a}^{2}}{q_{\star}^{2}} - 1 \right]$$
$$\frac{\beta_{t} q_{a}^{2}}{\epsilon} = \left[ \frac{q_{a}^{2}}{q_{\star}^{2}} \left( \frac{q_{a}^{2}}{q_{\star}^{2}} - 1 \right) \right]^{1/2}$$

# **Plot the Results**



2.  $I \propto 1/q_{\star}$ 



As  $v_*$  increases, *I* increases. This helps to prevent the separatrix from moving onto the plasma surface since less vertical field is required to maintain toroidal force balance.





Less vertical field is required. The separatrix stays away from the plasma surface.



No equilibrium limit. The separatrix does not move onto the plasma surface.

5.  $\beta_t$  vs.  $1/q_*$ 



# Summary

- 1. General HBT: covers all permissable  $\beta_t/\epsilon$ ,  $q_\star$  space
- 2. HBT at fixed *I*: exhibits an equilibrium limit
- 3. FCT at fixed  $q_a$ : no equilibrium limit