

Today

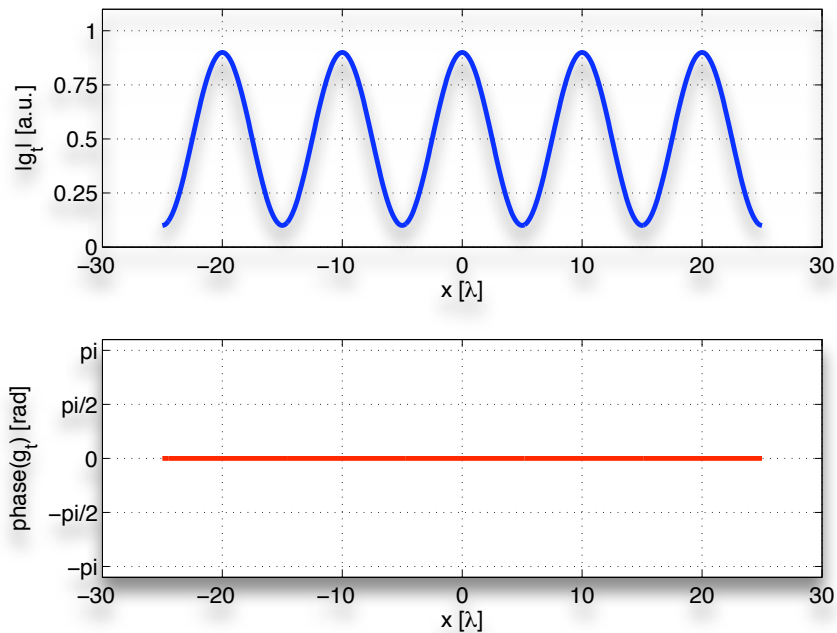
- Gratings:
 - Sinusoidal amplitude grating
 - Sinusoidal phase grating
 - in general: spatially periodic thin transparency

Wednesday

- Fraunhofer diffraction
- Fraunhofer patterns of typical apertures
- Spatial frequencies and Fourier transforms

Gratings

Amplitude grating



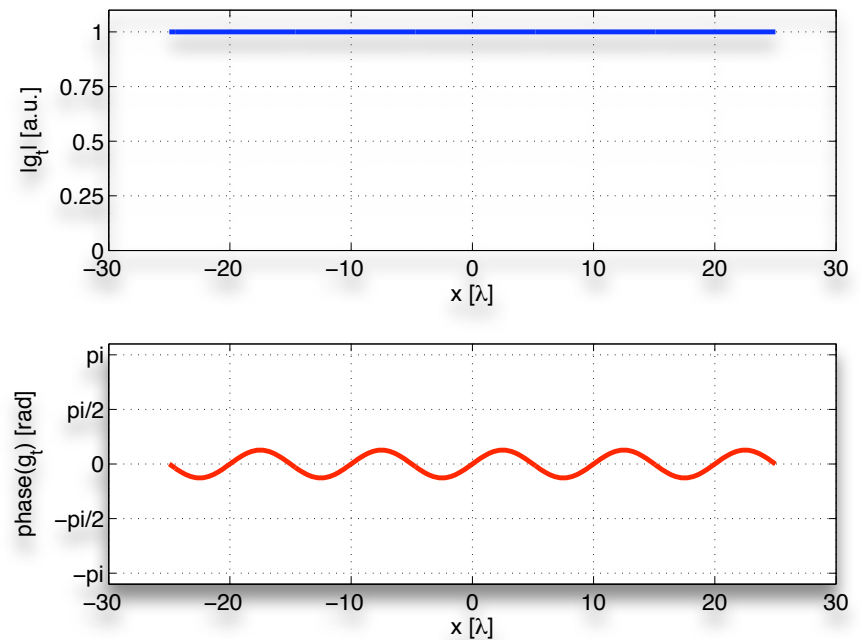
$$g_t(x) = \frac{1}{2} \left[1 + m \cos \left(2\pi \frac{x}{\Lambda} + \phi \right) \right]$$

$$|g_t(x)| = \frac{1}{2} \left[1 + m \cos \left(2\pi \frac{x}{\Lambda} + \phi \right) \right]$$

$$\angle g_t(x) = 0.$$

Λ : period; m : contrast; ϕ : phase shift

Phase grating



$$g_t(x) = \exp \left[i \frac{m}{2} \sin \left(2\pi \frac{x}{\Lambda} + \phi \right) \right]$$

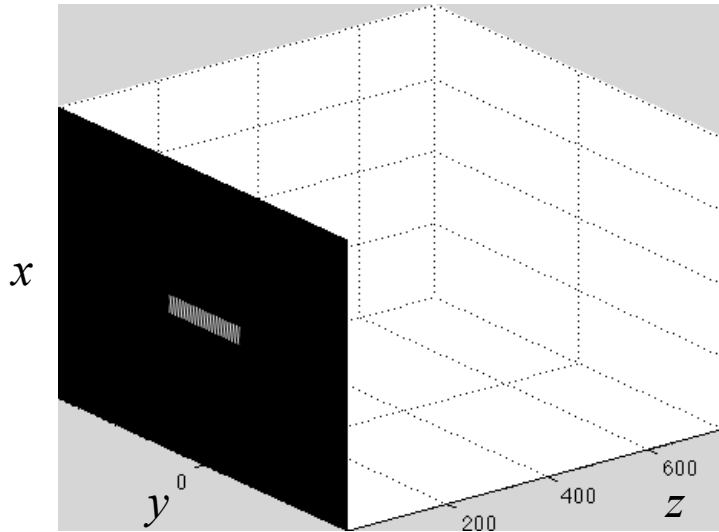
$$|g_t(x)| = 1$$

$$\angle g_t(x) = \frac{m}{2} \sin \left(2\pi \frac{x}{\Lambda} + \phi \right).$$

Λ : period; m : phase contrast; ϕ : phase shift

Gratings

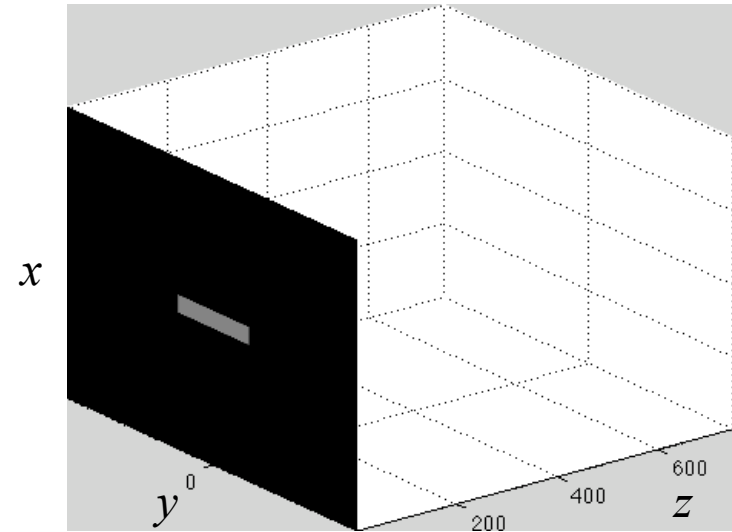
Amplitude grating



$$g_t(x) = \frac{1}{2} \left[1 + m \cos \left(2\pi \frac{x}{\Lambda} + \phi \right) \right]$$

$$\begin{aligned} \Lambda &= 2\lambda \\ m &= 1.0 \\ \phi &= 0. \end{aligned}$$

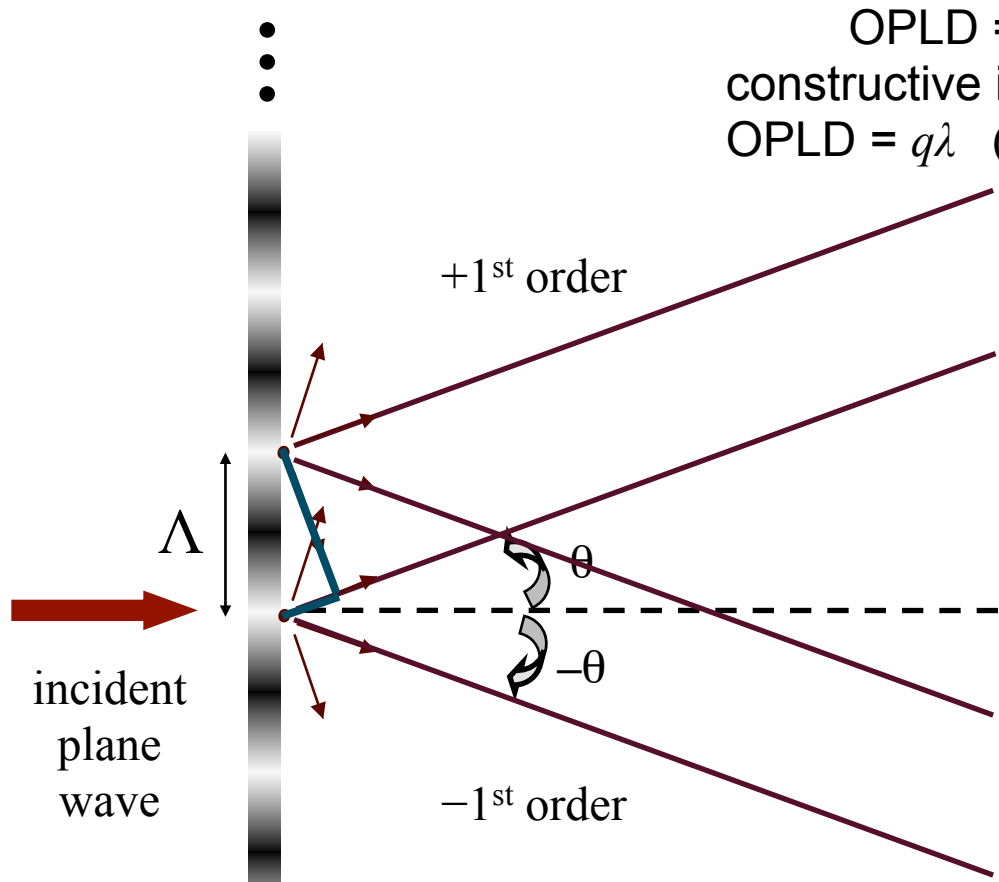
Phase grating



$$g_t(x) = \exp \left[i \frac{m}{2} \sin \left(2\pi \frac{x}{\Lambda} + \phi \right) \right]$$

$$\begin{aligned} \Lambda &= 2\lambda \\ m &= 8.44 \text{ rad} \\ \phi &= 0. \end{aligned}$$

Sinusoidal amplitude grating



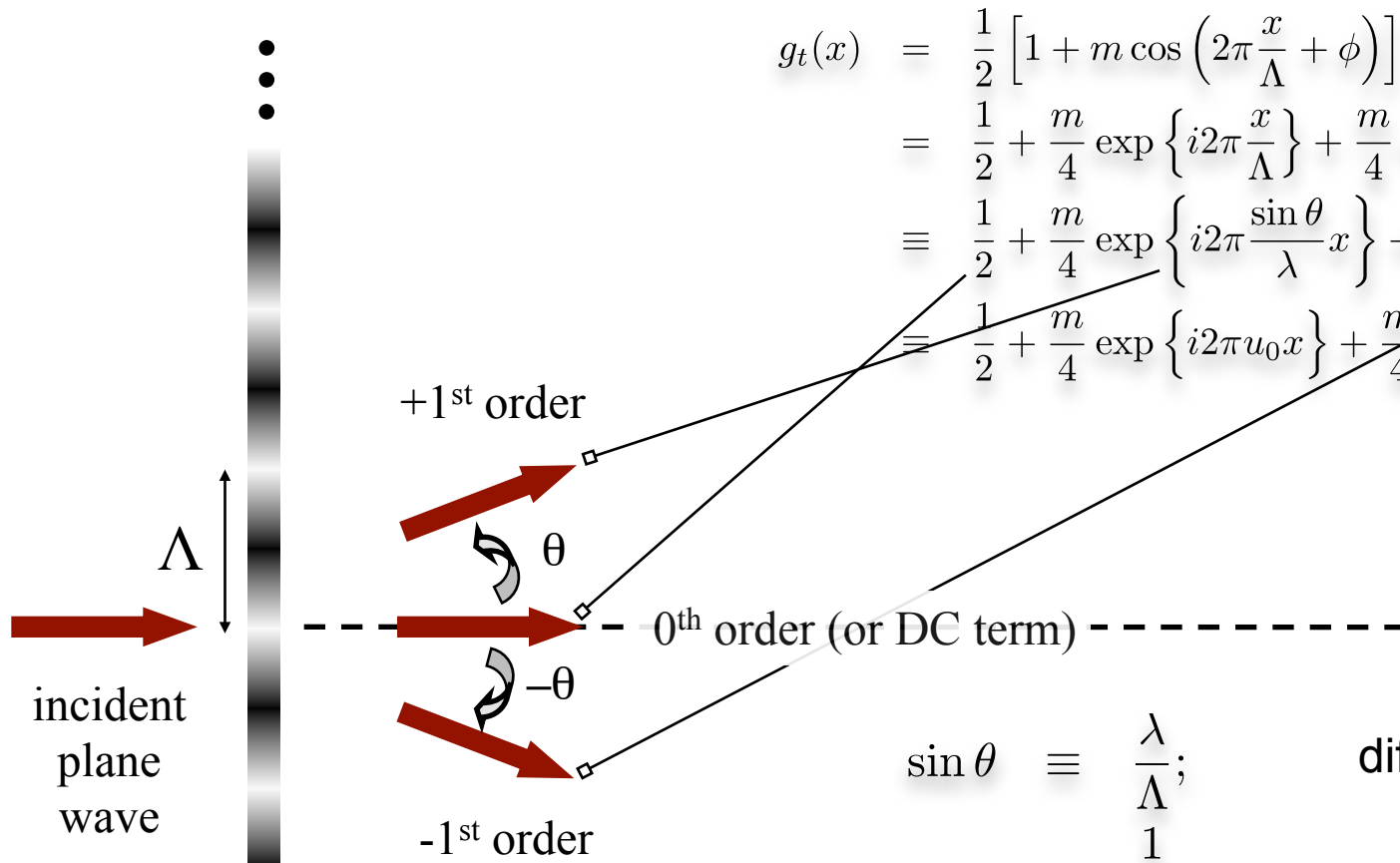
OPLD = $\Lambda \sin \theta$
 constructive interference if $\Rightarrow \sin \theta_q = \frac{\text{OPLD}}{\Lambda} = \frac{q\lambda}{\Lambda}$
 OPLD = $q\lambda$ (q integer)
 $\Rightarrow \theta_{\pm 1} \approx \pm \frac{\lambda}{\Lambda}$

Only the $n=0^{\text{th}}$ and $n=\pm 1^{\text{st}}$ diffraction orders are generated

The $n=0^{\text{th}}$ diffraction order is also known as the DC term

Grating period $\Lambda \equiv \frac{1}{u_0}$ u_0 : spatial frequency

Sinusoidal amplitude grating

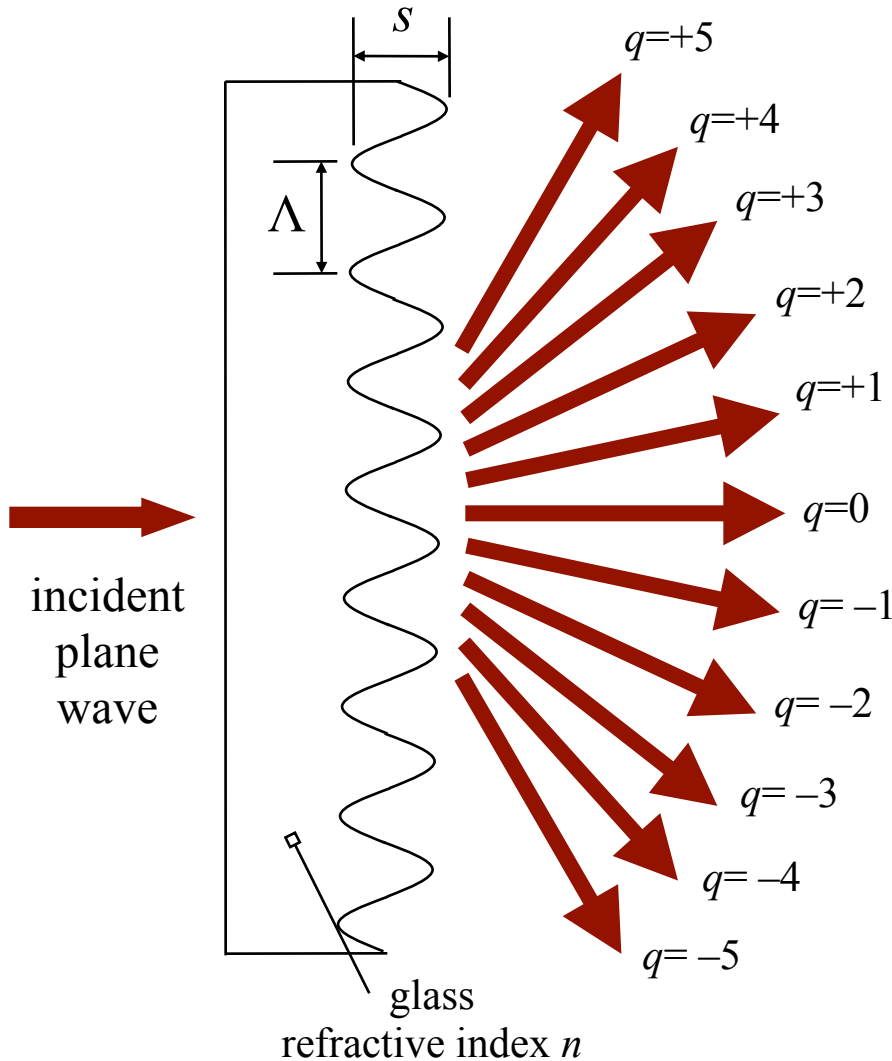


$$\begin{aligned}
 g_t(x) &= \frac{1}{2} \left[1 + m \cos \left(2\pi \frac{x}{\Lambda} + \phi \right) \right] \\
 &= \frac{1}{2} + \frac{m}{4} \exp \left\{ i2\pi \frac{x}{\Lambda} \right\} + \frac{m}{4} \exp \left\{ -i2\pi \frac{x}{\Lambda} \right\} \\
 &\equiv \frac{1}{2} + \frac{m}{4} \exp \left\{ i2\pi \frac{\sin \theta}{\lambda} x \right\} + \frac{m}{4} \exp \left\{ -i2\pi \frac{\sin \theta}{\lambda} x \right\} \\
 &\equiv \frac{1}{2} + \frac{m}{4} \exp \left\{ i2\pi u_0 x \right\} + \frac{m}{4} \exp \left\{ -i2\pi u_0 x \right\}.
 \end{aligned}$$

$$\begin{aligned}
 \sin \theta &\equiv \frac{\lambda}{\Lambda}; && \text{diffraction angle} \\
 u_0 &\equiv \frac{1}{\Lambda} \\
 &= \frac{\sin \theta}{\lambda}. && \text{spatial frequency}
 \end{aligned}$$

$$\eta_0 = \left(\frac{1}{2} \right)^2; \quad \eta_{\pm 1} = \left(\frac{m}{4} \right)^2 \quad \text{diffraction efficiencies}$$

Sinusoidal phase grating



“surface relief” grating

Transmission function

$$g_t(x) = \exp \left\{ i \frac{m}{2} \sin \left(2\pi \frac{x}{\Lambda} \right) \right\}$$

$$m \equiv 2\pi \frac{(n-1)s}{\lambda}$$

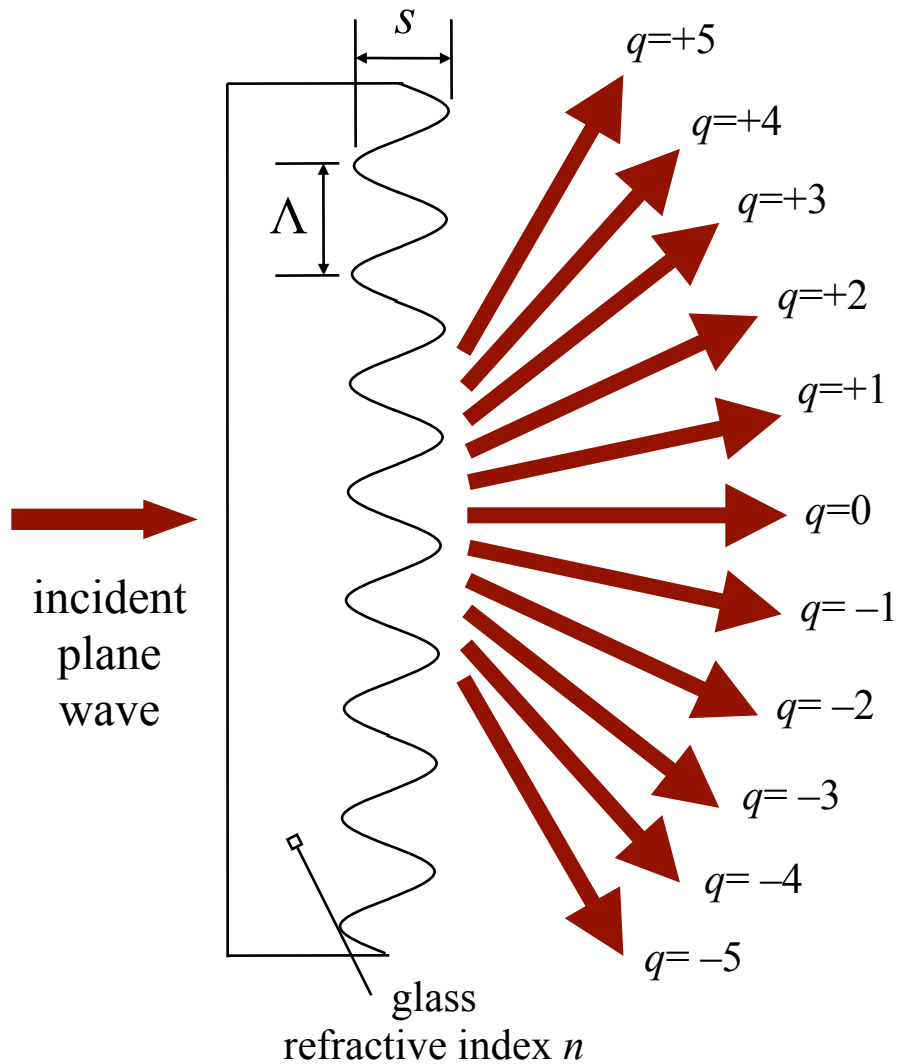
“phase contrast”

Useful math property:

$$\exp \{ i\alpha \sin(\theta) \} = \sum_{q=-\infty}^{+\infty} J_q(\alpha) \exp(iq\theta)$$

J_m : Bessel function of 1st kind,
 m -th order

Sinusoidal phase grating



Field after grating is expressed as:

$$g_+(x, z) = \exp \left\{ i \frac{m}{2} \sin \left(2\pi \frac{x}{\Lambda} \right) + i 2\pi \frac{z}{\lambda} \right\}$$

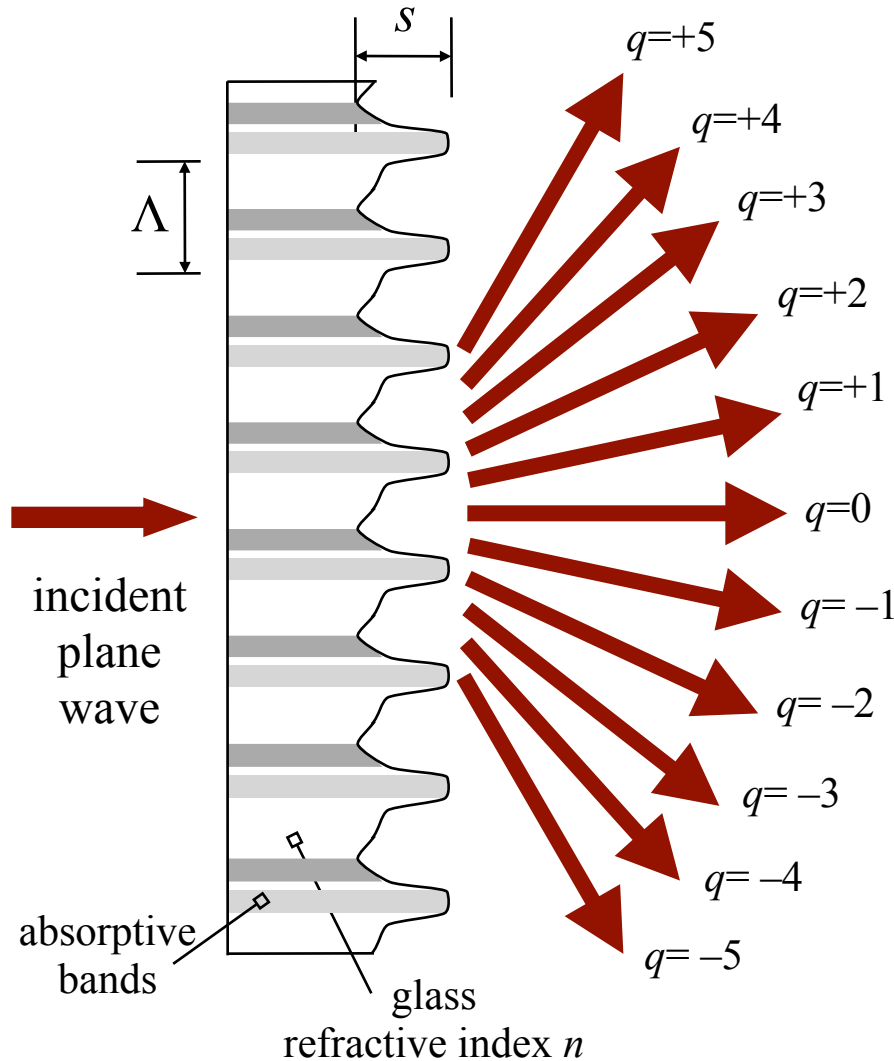
$$= \sum_{q=-\infty}^{+\infty} J_q \left(\frac{m}{2} \right) \times \text{amplitude}$$

$$\times \exp \left(i 2\pi q \frac{x}{\Lambda} + i 2\pi \frac{z}{\lambda} \right)$$

plane waves (diffraction orders),
 propagation angle
 $\sin \theta_q \approx \theta_q \equiv q \frac{\lambda}{\Lambda}$

The q -th exponential physically is
 a plane wave propagating at angle θ_q
i.e. the q -th diffraction order

Fresnel diffraction from a grating as a Fourier series



More generally, a periodic transparency's complex amplitude transmission may be expressed as a Fourier series expansion:

$$g_0(x) = \begin{cases} g_0(x), & 0 \leq x < \Lambda \\ 0, & \text{otherwise} \end{cases}$$

$$g_t(x) = g_0(x) \times \sum_{q=-\infty}^{+\infty} \delta(x - n\Lambda)$$

$$= \sum_{q=-\infty}^{+\infty} c_q \times$$

amplitudes

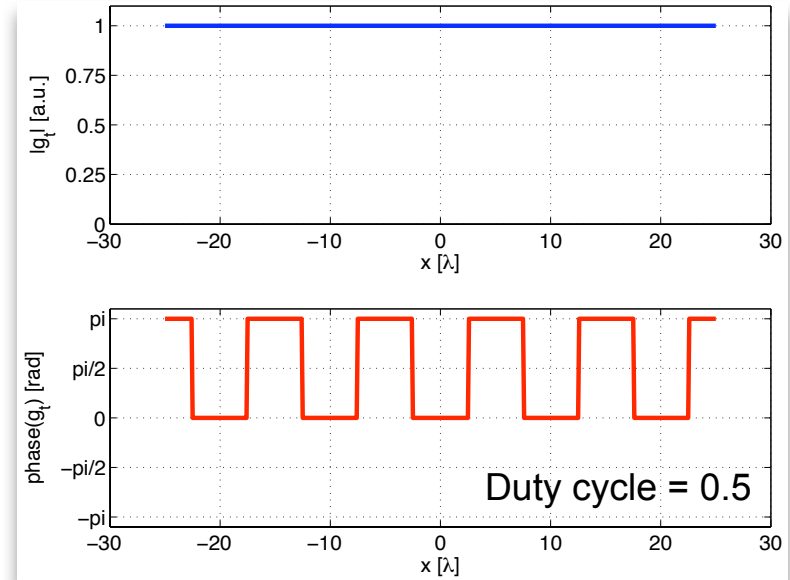
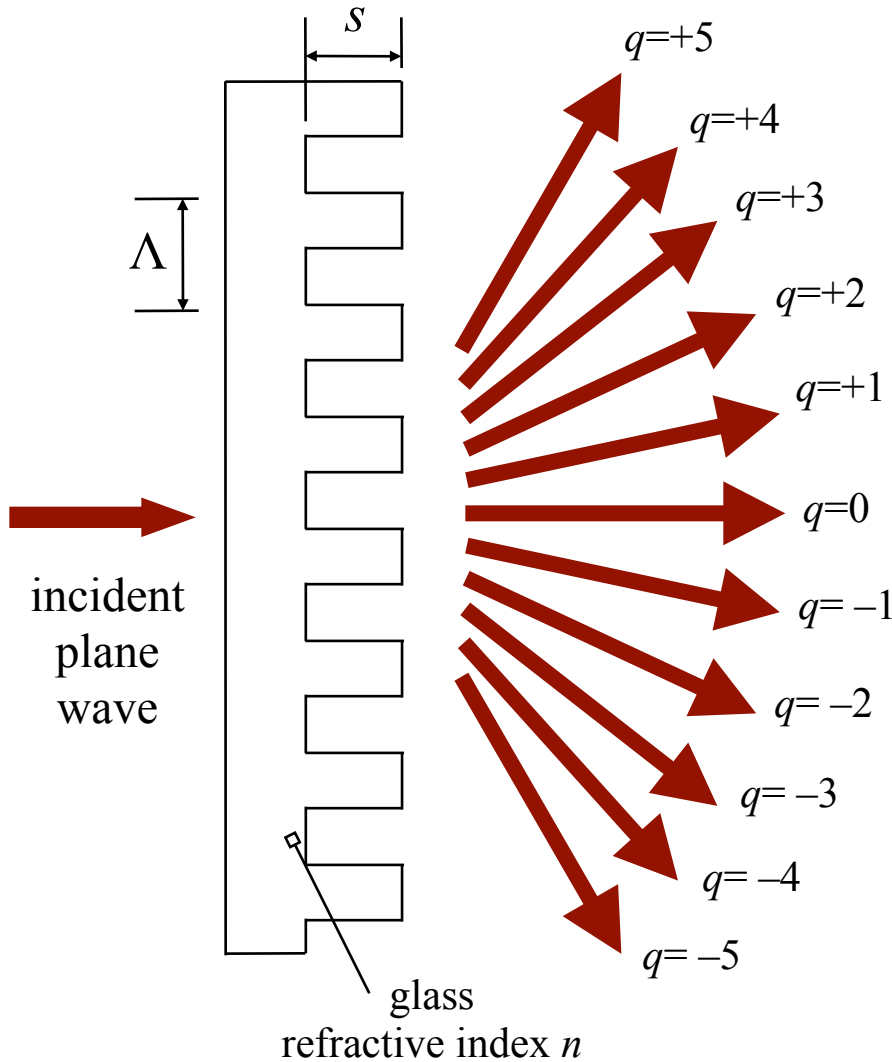
$$\times \exp\left(i2\pi q \frac{x}{\Lambda} + i2\pi \frac{z}{\lambda}\right)$$

plane waves (diffraction orders),
propagation angle

$$\sin \theta_q \approx \theta_q \equiv q \frac{\lambda}{\Lambda}$$

$$\eta_q = |c_q|^2 \quad \text{diffraction efficiencies}$$

Example: binary phase grating



$$g_0(x) = \begin{cases} 1, & 0 \leq |x| \leq \Lambda/4 \\ -1, & \Lambda/4 < |x| \leq \Lambda/2 \end{cases}$$

$$c_q = \frac{1}{\Lambda} \int_{-\Lambda/2}^{\Lambda/2} g_0(x) \exp\left\{i2\pi q \frac{x}{\Lambda}\right\} dx.$$

$$c_q = \text{sinc}\left(\frac{q}{2}\right) \quad \text{where} \quad \text{sinc}(\xi) \equiv \frac{\sin(\pi\xi)}{(\pi\xi)}.$$

$$\eta_{\pm q} = \left(\frac{2}{\pi q}\right)^2 \quad \text{for } q \text{ odd.}$$

$$\eta_{\pm 1} = \left(\frac{2}{\pi}\right)^2 \approx 40.53\%.$$

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2.71 / 2.710 Optics
Spring 2009

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