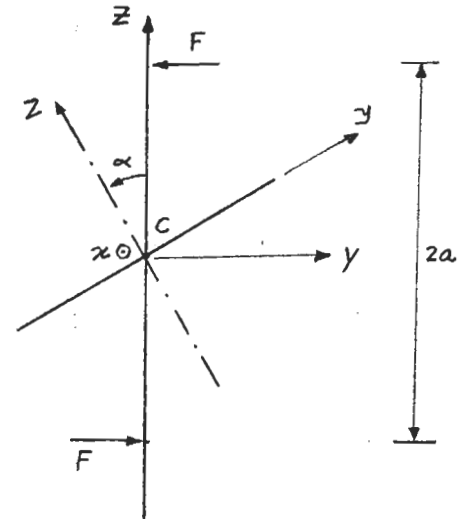
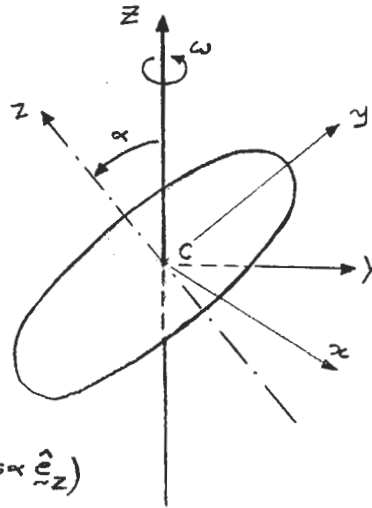


Problem 1

xyz coordinate system
is fixed to the disk
and rotates with ω
about Z axis:



(a)

$$\vec{\omega}_{\text{disk}} = \omega \hat{e}_Z = \omega (\sin\alpha \hat{e}_y + \cos\alpha \hat{e}_z)$$

$$\Rightarrow \begin{cases} \omega_x = 0 \\ \omega_y = \omega \sin\alpha \\ \omega_z = \omega \cos\alpha \end{cases}$$

$$I_x = I_y = \frac{1}{4} MR^2, \quad I_z = \frac{1}{2} MR^2 \quad \Rightarrow \quad [I]_C = \frac{1}{4} MR^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{bmatrix}$$

$$\vec{H}_C = [I]_C \vec{\omega} = \frac{1}{4} MR^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{bmatrix} \begin{Bmatrix} 0 \\ \omega \sin\alpha \\ \omega \cos\alpha \end{Bmatrix} = \frac{1}{4} MR^2 \omega (\sin\alpha \hat{e}_y + 2\cos\alpha \hat{e}_z)$$

Angular momentum of the disk about C

(b)

Since $\vec{v}_C = 0$, forces on the bearings are equal and in opposite directions:

$$\vec{\tau}_C = \frac{d\vec{H}_C}{dt} = \frac{1}{4} MR^2 \omega \left(\sin\alpha \frac{d\hat{e}_y}{dt} + 2\cos\alpha \frac{d\hat{e}_z}{dt} \right)$$

$\omega \hat{e}_z \times \hat{e}_y = -\omega \cos\alpha \hat{e}_x$
 $\omega \hat{e}_z \times \hat{e}_z = \omega \sin\alpha \hat{e}_x$

$$\therefore 2aF \hat{e}_x = \frac{1}{4} MR^2 \omega^2 \sin\alpha \cos\alpha \hat{e}_x$$

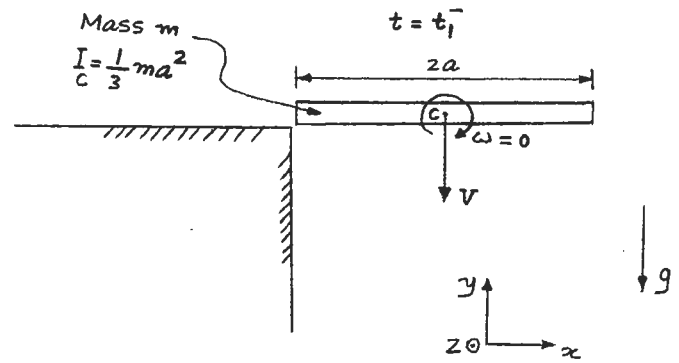
$$\Rightarrow F = \frac{MR^2 \omega^2 \sin 2\alpha}{16a} \quad \text{reaction forces at the bearings}$$

Note that force F rotates about Z axis with ω and is always in yz plane.

Problem 2

2

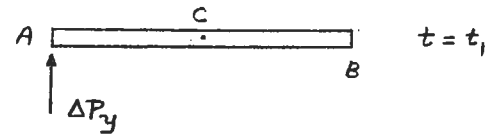
Assume the collision occurs at $t = t_1$:



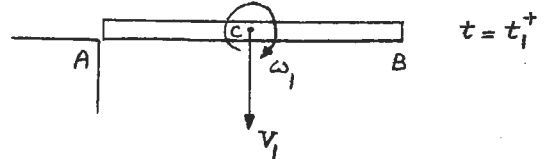
A vertical impulse acts on the end

A of the rod at $t = t_1$ (ΔP_y).

As a result, velocity of the center of mass would be v_1 and angular velocity ω_1 just after the impact.



(a) Energy is conserved in the



collision: $KE + PE|_{t=t_1^-} = KE + PE|_{t=t_1^+}$

Gravity does not have enough time to act:

$$PE|_{t=t_1^-} = PE|_{t=t_1^+}$$

$$KE = \frac{1}{2} m v_c^2 + \frac{1}{2} I_c \omega^2 \quad (I_c = \frac{1}{3} m a^2)$$

$$KE|_{t=t_1^-} = \frac{1}{2} m V^2,$$

$$KE|_{t=t_1^+} = \frac{1}{2} m v_1^2 + \frac{1}{6} m a^2 \omega_1^2$$

$$\therefore \underline{v_1^2 + \frac{1}{3} a^2 \omega_1^2 = V^2} \quad (1)$$

Angular momentum about point A on the table:

$$\tilde{\tau}_A = \frac{d}{dt} \tilde{H}_A + \tilde{v}_A \times \tilde{P}_A$$

Again, gravity does not have time to act. \rightarrow

$$\tilde{\tau}_A = 0 \quad \rightarrow \quad \frac{d}{dt} \tilde{H}_A = 0$$

$$\Rightarrow \tilde{H}_A|_{t=t_1^-} = \tilde{H}_A|_{t=t_1^+} \quad (\tilde{H}_A = \tilde{H}_C + \tilde{AC} \times m \tilde{v}_C)$$

$$H_A|_{t=t_1^-} = -m a V \hat{e}_z,$$

$$H_A|_{t=t_1^+} = -m a v_1 \hat{e}_z - \frac{1}{3} m a^2 \omega_1 \hat{e}_z$$

Problem 2

3

$$\therefore \underline{v_1 + \frac{1}{3} a \omega_1 = V} \quad (2)$$

$$(1), (2) \Rightarrow \left\{ \begin{array}{l} \cancel{v_1 = V}, \quad \cancel{\omega_1 = 0} \\ v_1 = \frac{V}{2}, \quad \omega_1 = \frac{3V}{2a} \end{array} \right. \quad \begin{array}{l} \text{angular velocity of the rod} \\ \text{just after the impact.} \end{array}$$

(b)

$$\begin{aligned} \underline{\underline{\tilde{v}_{A \text{ rod}} \Big|_{t=t_1^+} = \tilde{v}_C \Big|_{t=t_1^+} + \tilde{\omega}_{\text{rod}} \Big|_{t=t_1^+} \times \underline{CA}}} \\ = -v_1 \hat{e}_y + (-\omega_1 \hat{e}_z) \times (-a \hat{e}_x) \\ = -\frac{V}{2} \hat{e}_y + \left(-\frac{3V}{2a} \hat{e}_z\right) \times (-a \hat{e}_x) = \left(-\frac{V}{2} + \frac{3V}{2}\right) \hat{e}_y = \underline{\underline{V \hat{e}_y}} \end{aligned}$$

Velocity of the end of the rod that touched the table, just after the impact.

$$\underline{\underline{\tilde{v}_{A \text{ rod}} \Big|_{t=t_1^-} = -V \hat{e}_y}} \quad \text{just before the impact.}$$

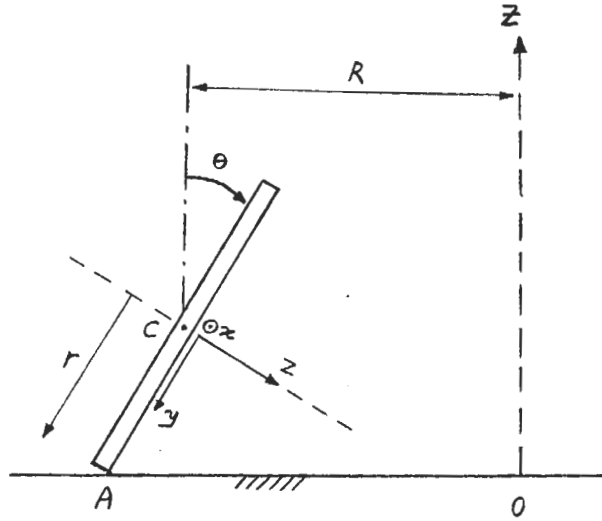
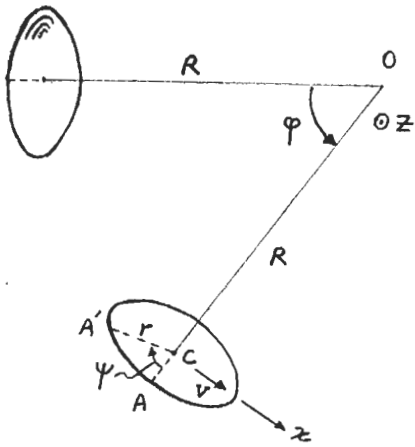
The result seems reasonable. Magnitude of the velocity of A on the rod is conserved

and $\tilde{v}_{A \text{ rod}}$ just changes the direction in the collision. This shows an elastic collision

which is expected when energy is conserved.

Problem 3

Assume thickness of the disk is small relative to its radius r .



xyz coordinate system rotates about z axis.

$$\underline{v}_C = v \hat{e}_x = R \dot{\varphi} \hat{e}_x \quad \underline{\omega}_{disk} = \dot{\varphi} \hat{e}_z + \dot{\psi} \hat{e}_z = \dot{\varphi} \left(-\cos\theta \hat{e}_y - \sin\theta \hat{e}_z \right) + \dot{\psi} \hat{e}_z$$

No slip. $\rightarrow \underline{v}_{A \text{ disk}} = \underline{0}$

$$\begin{aligned} \underline{v}_{A \text{ disk}} &= \underline{v}_C + \underline{\omega}_{disk} \times \underline{CA} \\ &= v \hat{e}_x + \left[-\frac{v}{R} \cos\theta \hat{e}_y + \left(\dot{\psi} - \frac{v}{R} \sin\theta \right) \hat{e}_z \right] \times (r \hat{e}_y) \\ &= \left[v - \left(\dot{\psi} - \frac{v}{R} \sin\theta \right) r \right] \hat{e}_x = \underline{0} \end{aligned}$$

$$\therefore \underline{\dot{\psi}} = v \left(\frac{1}{r} + \frac{1}{R} \sin\theta \right)$$

$$\Rightarrow \underline{\omega}_{disk} = -\frac{v}{R} \cos\theta \hat{e}_y + \frac{v}{r} \hat{e}_z$$

Angular velocity of the disk