

$$\underline{M} \ddot{x} + (\underline{C} + \underline{C}) \dot{x} + \underline{K} x = 0$$

$$x = q = \dot{q} =$$

Damped Oscillations (small) (about $q = q_0$; $x = q - q_0$)

(Holonomic scleronomous System)

$$(1) \quad \underline{M} \ddot{x} + \underline{C} \dot{x} + \underline{K} x = 0$$

$$\underline{M} = \underline{M}^T \text{ pos. def.}$$

$\underline{K} = \underline{K}^T \rightarrow$ pos. def. \rightarrow because system is Conservative

"damping nature"

If \int multiply from the left by \dot{x}^T

$$\Rightarrow \frac{d}{dt} \left[\frac{1}{2} \dot{x}^T \underline{M} \dot{x} + \frac{1}{2} x^T \underline{K} x \right] = - \dot{x}^T \underline{C} \dot{x} \Rightarrow \underline{C} \text{ is positive semi-definite}$$

quadratic terms in total energy

$$\underline{C} = \frac{1}{2} (\underline{C} + \underline{C}^T) + \frac{1}{2} (\underline{C} - \underline{C}^T)$$

Sym. skew-sym.

$$= - \dot{x}^T \left(\frac{1}{2} (\underline{C} + \underline{C}^T) \right) \dot{x}$$

only the symmetric part contribute to energy dissipation
if \underline{C} is positive definite then energy is decreased as long as the velocity is non-zero

Solution of (1) are of the form: $\ddot{x}(t) = a e^{\lambda t}$, $\lambda \in \mathbb{C}$

$$= a^{(\text{Re } \lambda + i \text{Im } \lambda)t} = a^{(\text{Re } \lambda)t} \left[\cos \text{Im } \lambda t + i \sin \text{Im } \lambda t \right]$$

\Rightarrow Substitution into (1):

$$(\lambda^2 \underline{M} + \lambda \underline{C} + \underline{K}) a = 0 \quad (a \neq 0)$$

$\rightarrow \det(\lambda^2 \underline{M} + \lambda \underline{C} + \underline{K}) = 0$ characteristic equation for λ

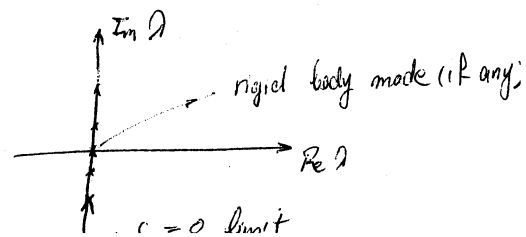
with $2n$ roots (real or Complex Conjugate pairs)

Roots on the Complex plane

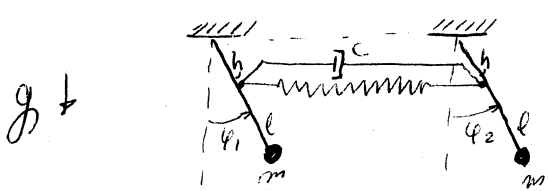
ω_+ : undamped natural frequency (for $\underline{C} = \underline{0}$)

$$\text{Re } \lambda_1 < 0$$

$$c = 0 \text{ limit}$$



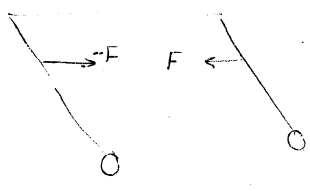
Example: Damped Pendulum-Spring System



c: damping coeff. for dashpot

we have seen: $M = \begin{pmatrix} ml^2 & 0 \\ 0 & ml^2 \end{pmatrix}$; $K = \begin{pmatrix} mgl + kh^2 & -kh^2 \\ -kh^2 & mgl + kh^2 \end{pmatrix}$

To find \underline{c} , first identify the generalized (non-potential) force in this system. active



$$E = -c \frac{d}{dt} [h(\sin \phi_2 - \sin \phi_1) \underline{i} - h(\cos \phi_2 - \cos \phi_1) \underline{j}]$$

$$= -ch \left[(\cos \phi_2 \dot{\phi}_2 - \cos \phi_1 \dot{\phi}_1) \underline{i} - (\sin \phi_2 \dot{\phi}_2 + \sin \phi_1 \dot{\phi}_1) \underline{j} \right]$$

$$\delta W_E = -F \cdot \delta r_1 + F \cdot \delta r_2 = \delta r_1 = \delta (h \sin \phi_1 - h \cos \phi_2)$$

$$= h \delta \phi_1 (\cos \phi_1 \underline{i} + \sin \phi_1 \underline{j})$$

$$\delta r_2 = h \delta \phi_2 (\cos \phi_2 \underline{i} + \sin \phi_2 \underline{j})$$

$$\delta W_E = ch^2 \delta \phi_1 [\cos(\phi_2 - \phi_1) \phi_2 - \phi_1] - ch^2 \delta \phi_2 [\phi_2 - \cos(\phi_2 - \phi_1) \phi_1]$$

$$\phi_1 = \alpha_1, \phi_2 = \alpha_2$$

$$\delta W_E = \underbrace{\delta \alpha_1 [-ch^2(\alpha_1 - \alpha_2) + O(2)]}_{C_{\alpha_1}} + \delta \alpha_2 \underbrace{[-ch^2(\alpha_2 - \alpha_1) + O(2)]}_{C_{\alpha_2}}$$

for linearized eq. of motion we need

$$\frac{\partial Q}{\partial \dot{\phi}} \Big|_{\phi=\phi_0=0} = \begin{pmatrix} -ch^2 & ch^2 \\ ch^2 & -ch^2 \end{pmatrix}$$

$\phi=0$

$$\underline{C} = \begin{pmatrix} ch^2 & -ch^2 \\ -ch^2 & ch^2 \end{pmatrix}; \text{ eig values } \lambda_1, \lambda_2$$

$$\frac{\partial Q}{\partial \dot{v}} \Big|_{\dot{v}=0} = \underline{-C}$$

Symmetric
Singular $\det \underline{C} = 0 \rightarrow \lambda_1, \lambda_2 = 0$
positive semi-definite
 $\text{tr}(\underline{C}) = 2ch^2 = \lambda_1 + \lambda_2$

because it is generalized force corresponding with damping
So \underline{C} is the negative of the value

$\Rightarrow \lambda_1 = 0, \lambda_2 > 0 \rightarrow \underline{C}$ pos. semi-definite \rightarrow acceptable

(makes sense physically these are small independent oscillations)

Forced Small Oscillations ($\epsilon = e$ for simplicity)

$$\underline{M} \ddot{\underline{x}} + \underline{k} \underline{x} = \underline{F}(t) = F \sin \omega t \quad (\text{Sinusoidal Forcing})$$

pass to modal (principal) coordinates

$$\underline{x} = \underline{\Phi} \underline{y}; \quad \underline{\Phi} = [\underline{a}_1, \dots, \underline{a}_n]$$

As earlier left multiplying (2) by $\underline{\Phi}^T$

$$\underline{\Phi}^T \underline{M} \underline{\Phi} \ddot{\underline{y}} + \underline{\Phi}^T \underline{k} \underline{\Phi} \underline{y} = \underline{\Phi}^T F \sin \omega t$$

$$\begin{pmatrix} m_1 & & 0 \\ & \ddots & \\ 0 & & m_n \end{pmatrix} \begin{pmatrix} k_1 & 0 \\ & \ddots & \\ 0 & & k_n \end{pmatrix} \underline{y} = \begin{Bmatrix} F_1 \\ \vdots \\ F_n \end{Bmatrix} \sin \omega t$$

$$\Rightarrow \ddot{y}_j + \frac{k_j}{m_j} y_j = \frac{F_j}{m_j} \sin \omega t; \quad j = 1, \dots, n$$

Solution: $y_j(t) = y_{j, \text{hom.}}(t) + y_{j, \text{par.}}(t)$

$$y_{j, \text{hom.}}(t) = C_1 \cos \omega_j t + C_2 \sin \omega_j t, \quad \omega_j^2 = \frac{k_j}{m_j}$$

Case (a) $\omega_j \neq \omega \Rightarrow y_{j, \text{par.}}(t) = A_j \sin \omega t \rightarrow A_j = \frac{F_{j,m}}{\omega_j^2 - \omega^2}$

$$= \frac{F_{j,m}}{\omega_j^2 - \omega^2} \sin \omega t$$

Case (b) $\omega_j = \omega$ (resonance) $\Rightarrow y_{j, \text{par.}}(t) = (A_j + B_j t) \sin \omega t + (C_j + D_j t) \cos \omega t$

$\Rightarrow \dots$ growing oscillations

Response Diagram

