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3.23 Electrical, Optical, and Magnetic Properties of Materials  
Fall 2007

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## 3.23 Fall 2007 – Lecture 17

# FERMAT'S FIRST THEOREM



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Pierre-Louis  
Moreau de  
Maupertuis



Hero

Abū 'Alī al-Ḥasan ibn  
al-Ḥasan ibn al-Haytham

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## Last time

1. Electric field, polarization, displacement, susceptibility
2. Maxwell's equations
3. Potentials and gauges
4. Electromagnetic waves (no free charges, currents)
5. Refractive index, phase and group velocity

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## Study

- (mostly read) Fox, Optical Properties of Solids: 1.1 to 1.4, 2.1 to 2.2.3, 3.1 to 3.3

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## Polarization, transversality of EM fields

$$\vec{E}(x, y, z, t) = \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$$

$\vec{\nabla} \cdot \vec{E} = 0$        $\vec{\nabla} \cdot \vec{H} = 0$        $\vec{k} = k_x \hat{x} + k_y \hat{y} + k_z \hat{z}$

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$
$$E_{0x} (-ik_x) e^{i(\omega t - \vec{k} \cdot \vec{r})}$$
$$E_{0x} k_x + E_{0y} k_y + E_{0z} k_z = 0 \quad \Rightarrow \quad \vec{E}_0 \cdot \vec{k} = 0$$

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## Boundary conditions (Gauss theorem)

$$\int_{\text{volume}} \vec{\nabla} \cdot \vec{B} dv = \int_{\text{surface}} \vec{B} \cdot \hat{n} dS = 0$$
$$\int_{\text{volume}} \vec{\nabla} \cdot \vec{D} dv = \int_{\text{surface}} \vec{D} \cdot \hat{n} dS = 4\pi \int_{\text{volume}} \rho dv$$

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## Boundary conditions

$$\hat{n} \cdot (\vec{B}_2 - \vec{B}_1) = 0$$

$$\hat{n} \cdot (\vec{D}_2 - \vec{D}_1) = \sigma \quad (\sigma = \text{surface charge density})$$

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## Boundary conditions (Stokes theorem)

$$\int_{\text{surface}} \vec{\nabla} \times \vec{E} \cdot \hat{n} dS = \int_{\text{line}} \vec{E} \cdot d\vec{r}$$

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## Boundary conditions

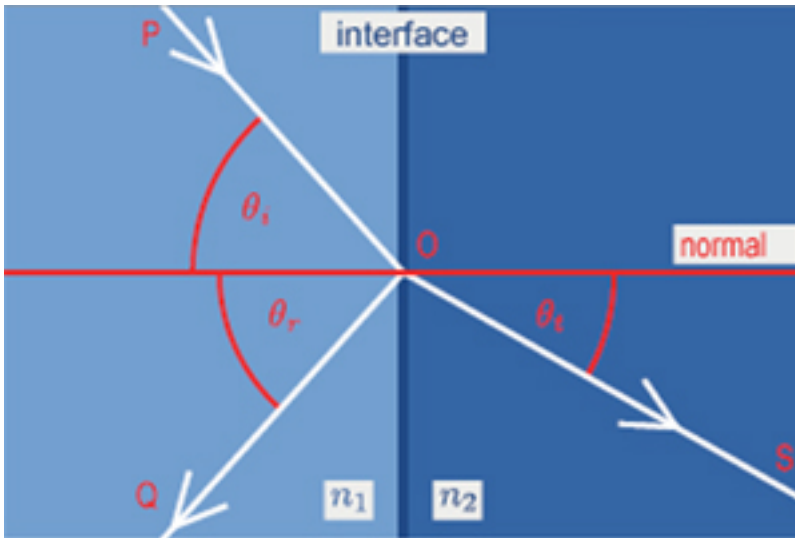
$$\hat{n} \times (\vec{E}_2 - \vec{E}_1) = 0$$

$$\hat{n} \times (\vec{H}_2 - \vec{H}_1) = \vec{K}$$

$$(\vec{K} = \text{surface current density})$$

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## Snell's law



$$\vec{E}_i e^{i\omega t - i\vec{k}_i \cdot \vec{r}}$$

incident wave

$$\vec{E}_r e^{i\omega t - i\vec{k}_r \cdot \vec{r}}$$

reflected wave

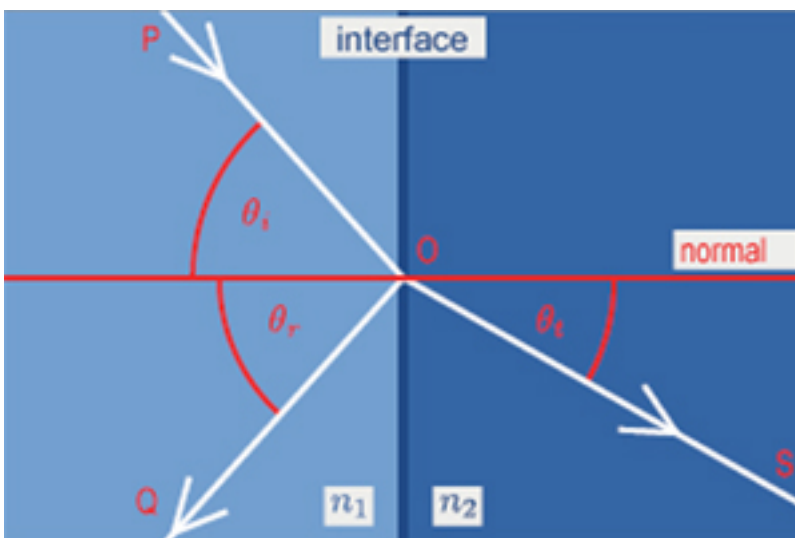
$$\vec{E}_t e^{i\omega t - i\vec{k}_t \cdot \vec{r}}$$

transmitted wave

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## Snell's law



$$|\vec{k}_i| = |\vec{k}_r| = \frac{\omega n_1}{c}$$

$$|\vec{k}_t| = \frac{\omega n_2}{c}$$

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## Snell's law

$$\left( \vec{k}_1 \cdot \vec{r} \right)_{x=0} = \left( \vec{k}'_1 \cdot \vec{r} \right)_{x=0} = \left( \vec{k}_2 \cdot \vec{r} \right)_{x=0}$$

$$(k_{1y}y + k_{1z}z) = (k'_{1y}y + k'_{1z}z) = (k_{2y}y + k_{2z}z) \rightarrow k_{1y} = k'_{1y} = k_{2y}$$

$$\text{and } k_{1z} = k'_{1z} = k_{2z}$$

$$\left( \vec{k}_{1t} \cdot \vec{r}_t \right) = \left( \vec{k}'_{1t} \cdot \vec{r}_t \right) = \left( \vec{k}_{2t} \cdot \vec{r}_t \right)$$

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## Snell's law

$$|\vec{k}_1| = |\vec{k}'_1| = n_1 \frac{\omega}{c}$$

$$|\vec{k}_2| = n_2 \frac{\omega}{c}$$

$$k_{iz} = k_{tz} \rightarrow |k_i| \sin \theta_1 = |k_t| \sin \theta_2$$

$$\frac{\omega n_1}{c} \sin \theta_1 = \frac{\omega n_2}{c} \sin \theta_2$$

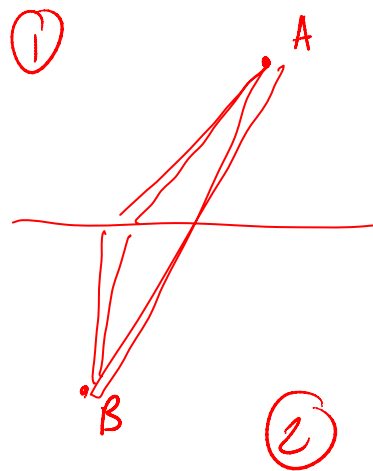
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## Snell's law

$$\left. \begin{aligned} k_{1z} &= |\vec{k}_1| \sin \theta_1 = n_1 \frac{\omega}{c} \sin \theta_1 \\ k_{2z} &= |\vec{k}_2| \sin \theta_2 = n_2 \frac{\omega}{c} \sin \theta_2 \end{aligned} \right\} n_1 \sin \theta_1 = n_2 \sin \theta_2$$

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## Principle of least action



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## Energy law

$$\vec{E} \cdot \vec{\nabla} \times \vec{H} - \vec{H} \cdot \vec{\nabla} \times \vec{E} = \frac{4\pi}{c} \vec{J} \cdot \vec{E} + \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} + \frac{1}{c} \vec{H} \cdot \frac{\partial \vec{B}}{\partial t}$$

$$\vec{E} \cdot \vec{\nabla} \times \vec{H} - \vec{H} \cdot \vec{\nabla} \times \vec{E} = -\vec{\nabla} \cdot (\vec{E} \times \vec{H})$$

$$\rightarrow \frac{4\pi}{c} \vec{J} \cdot \vec{E} + \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} + \frac{1}{c} \vec{H} \cdot \frac{\partial \vec{B}}{\partial t} + \vec{\nabla} \cdot (\vec{E} \times \vec{H}) = 0$$

Apply Gauss's theorem

$$\int_V \frac{4\pi}{c} \vec{J} \cdot \vec{E} dv + \int_V \left( \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} + \frac{1}{c} \vec{H} \cdot \frac{\partial \vec{B}}{\partial t} \right) dv + \int_S (\vec{E} \times \vec{H}) \cdot \hat{n} dS = 0$$

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## Energy law

$$\frac{1}{4\pi} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} = \frac{1}{4\pi} \vec{E} \cdot \frac{\partial \epsilon \vec{E}}{\partial t} = \frac{1}{8\pi} \frac{\partial \epsilon \vec{E}^2}{\partial t} = \frac{1}{8\pi} \frac{\partial (\vec{E} \cdot \vec{D})}{\partial t}$$

$$\frac{1}{4\pi} \vec{H} \cdot \frac{\partial \vec{B}}{\partial t} = \frac{1}{8\pi} \frac{\partial (\vec{H} \cdot \vec{B})}{\partial t}$$

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## Energy conservation

$$\int \vec{J} \cdot \vec{E} dv + \frac{\partial}{\partial t} \int \underbrace{(\vec{E} \cdot \vec{D} + \vec{H} \cdot \vec{B})}_{\substack{\text{total energy stored in electrical} \\ \text{and magnetic field} \\ \text{per volume}}} dv + \int \underbrace{(\vec{E} \times \vec{H})}_{\substack{\text{energy surface} \\ \text{flux per unit area}}} \cdot \hat{n} dS = 0$$

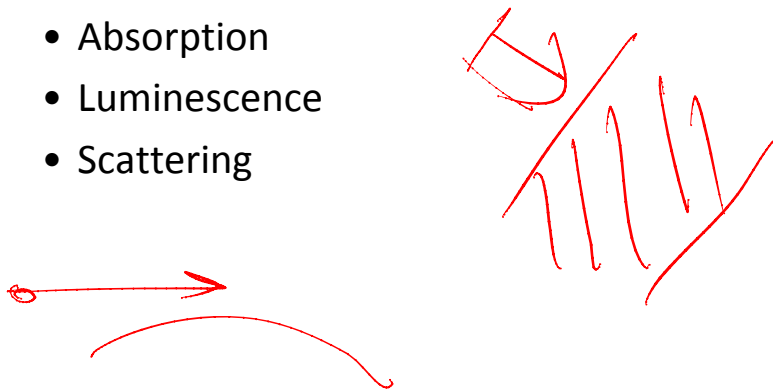
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$$\vec{S} = \frac{c}{4\pi} \vec{E} \times \vec{H}$$

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## Optical processes

- Reflection and refraction
- Absorption
- Luminescence
- Scattering



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## Optical coefficients

T: ratio of transmitted vs incident power

$R+T=1$  (no absorption, scattering)

Absorption:  $dI = -\alpha dz I \Rightarrow I(z) = I_0 e^{-\alpha z}$

Transmission:  $T = \frac{(1-R_1)e^{-\alpha l}}{(1-R_2)}$



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## Complex refractive index

$$\tilde{n} = n + ik$$

$$\vec{E}(x, y, z, t) = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$k = \frac{\omega n}{c} \Rightarrow \frac{\omega}{c} (n + ik)$$

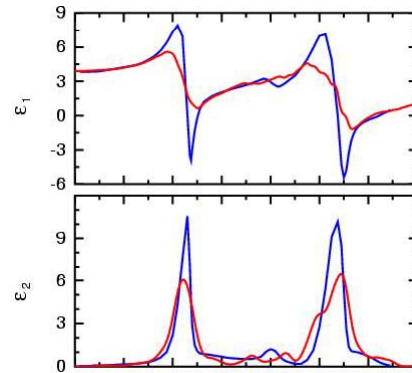
$$\vec{E}_0 e^{-kz} e^{i(\omega n z/c - \omega t)}$$

$$\epsilon = n^2 \Rightarrow \tilde{\epsilon} = \tilde{n}^2 = \epsilon_1 + i\epsilon_2$$

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# Complex refractive index

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 Please see any image of the structure of amorphous silica,  
 such as <http://www.research.ibm.com/amorphous/figure1.gif>.



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## Modeling Optical Constants with a Damped Harmonic Oscillator

$$\underbrace{m_0 \frac{d^2 X}{dt^2}}_{\text{acceleration}} + \underbrace{m_0 \gamma \frac{dX}{dt}}_{\text{dissipation}} + \underbrace{m_0 \omega_0^2 X}_{\text{harmonic restoring force}} = \underbrace{-eE(t)}_{\text{time dependent electric field}}$$

$E(t) = E e^{-i\omega t - \phi}$   
 $X(t) = X_0 e^{-\omega t - \phi'}$   
 $-m_0 \omega^2 X_0 e^{-i\omega t} - i m_0 \gamma \omega X_0 e^{-i\omega t} + m_0 \omega_0^2 X_0 e^{-i\omega t} = -\epsilon_0 e^{-i\omega t}$

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## Modeling Optical Constants with a Damped Harmonic Oscillator

$$X_0 = \frac{-eE_0}{m_0(\omega_0^2 - \omega^2 - i\gamma\omega)}$$

$\chi = \chi_0 e^{-i\omega t}$

$$P_{\text{resonant}} = Np = -NeX = \underbrace{\frac{Ne^2}{m_0(\omega_0^2 - \omega^2 - i\gamma\omega)}}_{\alpha} E$$

$$D = E + 4\pi P + 4\pi P_{\text{resonant}} = E + 4\pi\chi E + 4\pi \underbrace{\frac{Ne^2}{m_0(\omega_0^2 - \omega^2 - i\gamma\omega)}}_{\alpha} E = \varepsilon E$$

Atomic polarizability =  $\alpha$

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## Modeling Optical Constants with a Damped Harmonic Oscillator

$$\varepsilon = 1 + 4\pi\chi + 4\pi \underbrace{\frac{Ne^2(\omega_0^2 - \omega^2)}{m_0((\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2)}}_{\varepsilon_1} - i4\pi \underbrace{\frac{Ne^2\gamma\omega}{m_0((\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2)}}_{\varepsilon_2}$$

$$\varepsilon = (n + ik)^2 = \underbrace{n^2 - k^2}_{\varepsilon_1} + i \underbrace{2nk}_{\varepsilon_2}$$

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# Amorphous silica

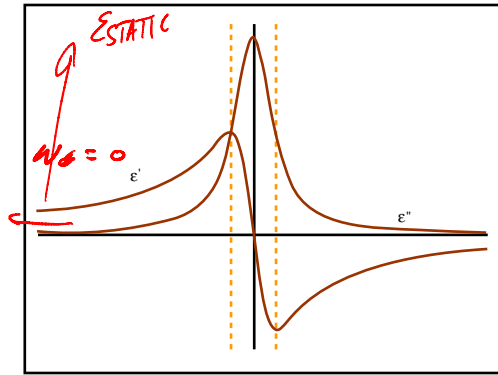
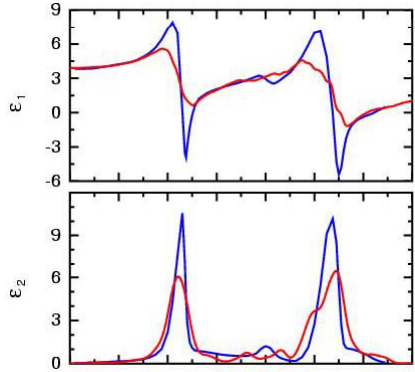


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# Optical materials

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Please see: Fig. 1.4 in Fox, Mark. *Optical Properties of Solids*. Oxford, England: Oxford University Press, 2001.

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# Infrared active modes

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"Ab initio Calculation of Phonon Dispersions in Semiconductors." *Physical Review B* 43 (March 15, 1991): 7231-7242.

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## Transition rate for direct absorption

$$W_{i \rightarrow f} = \frac{2\pi}{\hbar} |\langle f | H' | i \rangle|^2 \rho_f(\hbar\omega)$$

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such as [http://ecee.colorado.edu/~bart/book/book/chapter2/gif/fig2\\_3\\_6.gif](http://ecee.colorado.edu/~bart/book/book/chapter2/gif/fig2_3_6.gif).

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