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3.23 Electrical, Optical, and Magnetic Properties of Materials

Fall 2007

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3.23 Fall 2007 – Lecture 13

THE LAW OF MASS ACTION

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Last time

1. Band structure of oxides (perovskites), semiconductors (silicon, and compared with lead), late (fcc) transition metals (same period, or same group), graphene and nanotubes
2. Independent electron gas: states, energy, density, DOS
3. DOS of massive and massless bands in 1, 2 and 3 dimensions
4. Statistics of classical and quantum particles, Fermi-Dirac distribution, chemical potential

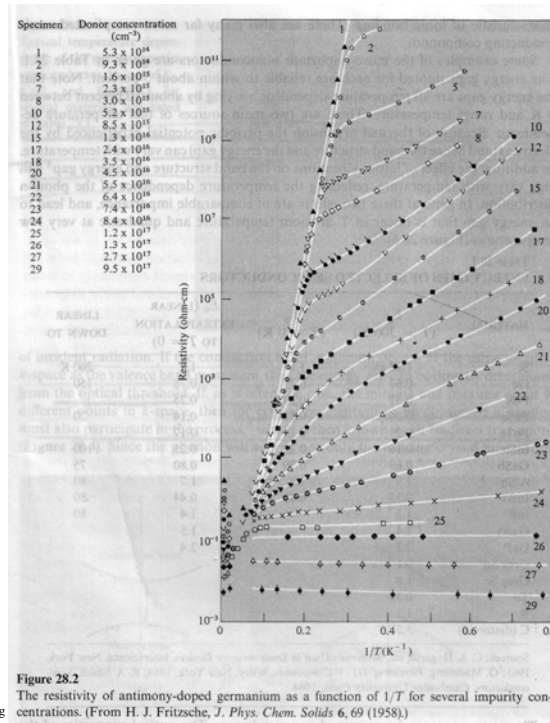
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Study

- Chap 6 Singleton,
or, much better,
- Chap 28 (Homogeneous semiconductors)
Ashcroft-Mermin (to be posted)

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Sb-doped Germanium



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Semiconductors

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Please see any graph of semiconductor band gaps vs. lattice constants, such as http://www.tf.uni-kiel.de/matwis/amat/semi_en/kap_5/illustr/bandgap_misfit.gif

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Valence+conduction bands in Si

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Please see: Fig. 6.1 in Singleton, John. *Band Theory and Electronic Properties of Solids*. Oxford, England: Oxford University Press, 2001.

Band structure of Si, Ge, GaAs

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Please see any image of Si, Ge, and GaAs energy bands,
such as http://ecee.colorado.edu/~bart/book/book/chapter2/gif/fig2_3_6.gif.

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Conduction band minima (in **3d**)

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Please see Fig. 19.9 in Marder, Michael P. *Condensed Matter Physics*. New York, NY: Wiley-Interscience, 2000.

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Optical absorption in Ge

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Please see: Fig. 6.3 and 6.4 in Singleton, John. *Band Theory and Electronic Properties of Solids*. Oxford, England: Oxford University Press, 2001.

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Impress your examiners (orals)

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Please see Table 19.1 in Marder, Michael P. *Condensed Matter Physics*. New York, NY: Wiley-Interscience, 2000.

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Number of carriers at thermal equilibrium

$$n_c(T) = \int_{\epsilon_c}^{\infty} d\epsilon g_c(\epsilon) \frac{1}{e^{(\epsilon-\mu)/k_B T} + 1}$$

$$p_v(T) = \int_{-\infty}^{\epsilon_v} d\epsilon g_v(\epsilon) \left(1 - \frac{1}{e^{(\epsilon-\mu)/k_B T} + 1} \right)$$

$$= \int_{-\infty}^{\epsilon_v} d\epsilon g_v(\epsilon) \left(\frac{1}{e^{(\mu-\epsilon)/k_B T} + 1} \right)$$

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Please see Fig. 2.16 in Pierret, Robert F. *Semiconductor Device Fundamentals*. Reading, MA: Addison-Wesley, 1996.

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Conduction and valence DOS (non-degenerate sc, isotropic effective mass)

$$\frac{1}{e^{(\epsilon-\mu)/k_B T} + 1} \approx e^{-(\epsilon-\mu)/k_B T}, \quad \epsilon > \epsilon_c$$

$$\frac{1}{e^{(\mu-\epsilon)/k_B T} + 1} \approx e^{-(\mu-\epsilon)/k_B T}, \quad \epsilon < \epsilon_v$$

$$n_c(T) = \underbrace{\int_{\epsilon_c}^{\infty} d\epsilon g_c(\epsilon)}_{N_c(T)} e^{-(\epsilon_c-\mu)/k_B T} e^{-(\mu-\epsilon_c)/k_B T}$$

$$p_v(T) = \underbrace{\int_{-\infty}^{\epsilon_v} d\epsilon g_v(\epsilon)}_{P_v(T)} e^{-(\epsilon_v-\mu)/k_B T} e^{-(\mu-\epsilon_v)/k_B T}$$

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Please see: Fig. 18 in Kittel, Charles. "Introduction to Solid State Physics." Chapter 8 in *Semiconductor Crystals*. New York, NY: John Wiley & Sons, 2004.

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Density of available states

$$\underbrace{\int_{\varepsilon_c}^{\infty} d\varepsilon g_c(\varepsilon) e^{-(\varepsilon - \varepsilon_c)/k_B T}}_{N_c(T)} \quad \begin{aligned} g_c(\varepsilon) &= \sqrt{2(\varepsilon - \varepsilon_c)} \frac{m_c^{3/2}}{\pi^2 \hbar^3} \\ g_v(\varepsilon) &= \sqrt{2(\varepsilon_v - \varepsilon)} \frac{m_v^{3/2}}{\pi^2 \hbar^3} \end{aligned}$$

$$N_c(T) = \frac{1}{4} \left(\frac{2m_c k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$P_v(T) = \frac{1}{4} \left(\frac{2m_v k_B T}{\pi \hbar^2} \right)^{3/2}$$

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Miracle ! Law of Mass Action

$$\begin{aligned} n_c(T) &= \underbrace{\int_{\varepsilon_c}^{\infty} d\varepsilon g_c(\varepsilon) e^{-(\varepsilon - \varepsilon_c)/k_B T}}_{N_c(T)} e^{-(\varepsilon_c - \mu)/k_B T} & N_c(T) &= 2.5 \left(\frac{m_c}{m} \right)^{3/2} \left(\frac{T}{300K} \right)^{3/2} 10^{19} / \text{cm}^3 \\ p_v(T) &= \underbrace{\int_{-\infty}^{\varepsilon_v} d\varepsilon g_v(\varepsilon) e^{-(\varepsilon_v - \varepsilon)/k_B T}}_{P_v(T)} e^{-(\mu - \varepsilon_v)/k_B T} & P_v(T) &= 2.5 \left(\frac{m_v}{m} \right)^{3/2} \left(\frac{T}{300K} \right)^{3/2} 10^{19} / \text{cm}^3 \end{aligned}$$

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Intrinsic case

$$n_c(T) = p_v(T) \equiv n_i(T)$$

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Intrinsic case

$$n_c(T) = \underbrace{\int_{\varepsilon_c}^{\infty} d\varepsilon g_c(\varepsilon) e^{-(\varepsilon - \varepsilon_c)/k_B T}}_{N_c(T)} e^{-(\varepsilon_c - \mu)/k_B T}$$
$$p_v(T) = \underbrace{\int_{-\infty}^{\varepsilon_v} d\varepsilon g_v(\varepsilon) e^{-(\varepsilon_v - \varepsilon)/k_B T}}_{P_v(T)} e^{-(\mu - \varepsilon_v)/k_B T}$$

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Extrinsic case

$$n_c(T) - p_v(T) = \Delta n$$

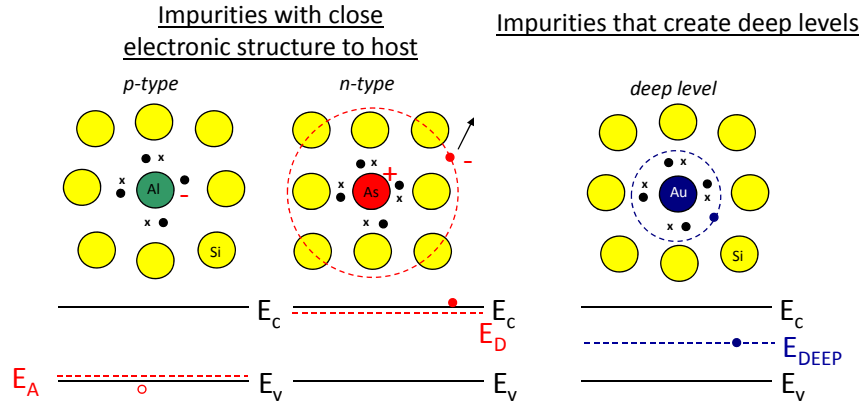
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Extrinsic case

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Impurity levels

- Adding impurities can lead to controlled domination of one carrier type
 - n-type is dominated by electrons
 - p-type is dominated by holes
- Adding other impurities can degrade electrical properties



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Impurity states as “embedded” hydrogen atoms

- Consider the weakly bound 5th electron in Phosphorus as a modified hydrogen atom
- For hydrogenic donors or acceptors, we can think of the electron or hole, respectively, as an orbiting electron around a net fixed charge
- We can estimate the energy to free the carrier into the conduction band or valence band by using a modified expression for the energy of an electron in the H atom

$$E_n = \frac{me^4}{8\epsilon_0^2 h^2 n^2} = -\frac{13.6}{n^2} \text{ (eV)}$$

$$E_n = \frac{me^4}{8\epsilon_0^2 h^2 n^2} \xrightarrow{\frac{e^2}{\epsilon_r} = e^2} \frac{m^* e^4}{8\epsilon_0^2 h^2 n^2 \epsilon_r^2} = -\frac{13.6}{n^2} \frac{m^*}{m} \frac{1}{\epsilon_r^2}$$

- Thus, for the ground state n=1, we can see already that since ϵ_r is on the order of 10, the binding energy of the carrier to the impurity atom is <0.1eV
- Expect that many carriers are then ionized at room T
 - B acceptor in Si: 0.046 eV
 - P donor in Si: 0.044 eV
 - As donor in Si: 0.049 eV

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Temperature dependence of majority carriers

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Please see Fig. 2.22 in Pierret, Robert F. *Semiconductor Device Fundamentals*. Reading, MA: Addison-Wesley, 1996.

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(Inverse temperature plot)

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Please see: Fig. 6.12 in Singleton, John. *Band Theory and Electronic Properties of Solids*. Oxford, England: Oxford University Press, 2001.

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